

# THE BASIC STRUCTURE OF SCATTERING AMPLITUDES



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(these slides are supposed to be integrated with the blackboard notes)

References:

- Ossola, & P.M., *On the Integrand Reduction Method for Two-Loop Scattering Amplitudes*, 1107.6041 [hep-ph]
- Badger, Frellersvig, Zhang, *Hepta-cuts of Two-Loop Scattering Amplitudes*, 1202.2019 [hep-ph]
- Zhang, *Integrand-level Reduction of Loop Amplitudes by Computational Algebraic Geometry Methods* 1205.5707 [hep-ph]
- Mirabella, Peraro, Ossola, & P.M., *Scattering Amplitudes from Multivariate Polynomial Division*, 1205.7087 [hep-ph]

# ONE-LOOP INTEGRAND DECOMPOSITION

Ossola, Papadopoulos, Pittau  
Ellis, Giele, Kunszt, Melnikov

$$\mathcal{A}_n^{\text{one-loop}} = \int d^{-2\epsilon}\mu \int d^4q A_n(q, \mu^2), \quad A_n(q, \mu^2) \equiv \frac{\mathcal{N}_n(q, \mu^2)}{\bar{D}_0 \bar{D}_1 \cdots \bar{D}_{n-1}} \quad \bar{D}_i = (\bar{q} + p_i)^2 - m_i^2 = (q + p_i)^2 - m_i^2 - \mu^2$$

We use a bar to denote objects living in  $d = 4 - 2\epsilon$  dimensions

$$\not{\bar{q}} = \not{q} + \not{\mu}, \quad \text{with} \quad \bar{q}^2 = q^2 - \mu^2.$$

$$\mathcal{A}_n^{\text{one-loop}} = c_{5,0} \text{ (pentagon) } + c_{4,0} \text{ (square) } + c_{4,4} \text{ (square with } d+4 \text{) } + c_{3,0} \text{ (triangle) } + c_{3,7} \text{ (triangle with } d+2 \text{) } + c_{2,0} \text{ (circle) } + c_{2,9} \text{ (circle with } d+2 \text{) } + c_{1,0} \text{ (circle) }$$

Passarino, Veltman; Tarasov

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Passarino, Veltman; Tarasov

## ☑ @ THE INTEGRAND-LEVEL

$$A_n(q, \mu^2) \neq \frac{c_{5,0}}{\bar{D}_0 \bar{D}_1 \bar{D}_2 \bar{D}_3 \bar{D}_4} + \frac{c_{4,0} + c_{4,4} \mu^4}{\bar{D}_0 \bar{D}_1 \bar{D}_2 \bar{D}_3} + \frac{c_{3,0} + c_{3,7} \mu^2}{\bar{D}_0 \bar{D}_1 \bar{D}_2} + \frac{c_{2,0} + c_{2,9} \mu^2}{\bar{D}_0 \bar{D}_1} + \frac{c_{1,0}}{\bar{D}_0}$$

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We use a bar to denote objects living in  $d = 4 - 2\epsilon$  dimensions  $\not{q} = \not{q} + \not{\mu}$ , with  $\bar{q}^2 = q^2 - \mu^2$ .

$$\mathcal{A}_n^{\text{one-loop}} = c_{5,0} \text{ (pentagon) } + c_{4,0} \text{ (square) } + c_{4,4} \text{ (square with } d+4 \text{) } + c_{3,0} \text{ (triangle) } + c_{3,7} \text{ (triangle with } d+2 \text{) } + c_{2,0} \text{ (circle) } + c_{2,9} \text{ (circle with } d+2 \text{) } + c_{1,0} \text{ (circle) }$$

Passarino, Veltman; Tarasov

## ☑ @ THE INTEGRAND-LEVEL

$$\begin{aligned} A_n(q, \mu^2) &\neq \frac{c_{5,0}}{\bar{D}_0 \bar{D}_1 \bar{D}_2 \bar{D}_3 \bar{D}_4} + \frac{c_{4,0} + c_{4,4} \mu^4}{\bar{D}_0 \bar{D}_1 \bar{D}_2 \bar{D}_3} + \frac{c_{3,0} + c_{3,7} \mu^2}{\bar{D}_0 \bar{D}_1 \bar{D}_2} + \frac{c_{2,0} + c_{2,9} \mu^2}{\bar{D}_0 \bar{D}_1} + \frac{c_{1,0}}{\bar{D}_0} \\ &= \frac{c_{5,0} + f_{01234}(q, \mu^2)}{\bar{D}_0 \bar{D}_1 \bar{D}_2 \bar{D}_3 \bar{D}_4} + \frac{c_{4,0} + c_{4,4} \mu^4 + f_{0123}(q, \mu^2)}{\bar{D}_0 \bar{D}_1 \bar{D}_2 \bar{D}_3} + \frac{c_{3,0} + c_{3,7} \mu^2 + f_{012}(q, \mu^2)}{\bar{D}_0 \bar{D}_1 \bar{D}_2} + \frac{c_{2,0} + c_{2,9} \mu^2 + f_{01}(q, \mu^2)}{\bar{D}_0 \bar{D}_1} + \frac{c_{1,0} + f_0(q, \mu^2)}{\bar{D}_0} \end{aligned}$$

☑ SPURIOUS TERMS  $\int d^{-2\epsilon} \mu \int d^4 q \frac{f_{i_1 i_2 \dots i_n}(q, \mu^2)}{\bar{D}_{i_1} \bar{D}_{i_2} \cdots \bar{D}_{i_n}} = 0.$

Ossola, Papadopoulos, Pittau

Parametric form of the residues: known

✓ MULTI-(PARTICLE)-POLE DECOMPOSITION

$$A(\bar{q}) = \sum_{i \ll m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_l \bar{D}_m} + \sum_{i \ll l}^{n-1} \frac{\Delta_{ijkl}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_l} + \sum_{i \ll k}^{n-1} \frac{\Delta_{ijk}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k} + \sum_{i < j}^{n-1} \frac{\Delta_{ij}(\bar{q})}{\bar{D}_i \bar{D}_j} + \sum_i^{n-1} \frac{\Delta_i(\bar{q})}{\bar{D}_i},$$

✓ INTEGRAND REDUCTION FORMULA

$$N(\bar{q}) = \sum_{i \ll m}^{n-1} \Delta_{ijklm}(\bar{q}) \prod_{h \neq i,j,k,l,m}^{n-1} \bar{D}_h + \sum_{i \ll l}^{n-1} \Delta_{ijkl}(\bar{q}) \prod_{h \neq i,j,k,l}^{n-1} \bar{D}_h + \sum_{i \ll k}^{n-1} \Delta_{ijk}(\bar{q}) \prod_{h \neq i,j,k}^{n-1} \bar{D}_h + \sum_{i < j}^{n-1} \Delta_{ij}(\bar{q}) \prod_{h \neq i,j}^{n-1} \bar{D}_h + \sum_i^{n-1} \Delta_i(\bar{q}) \prod_{h \neq i}^{n-1} \bar{D}_h,$$

Use your favourite generator,  
(for Feynman Diagrams, or for products of tree-amplitudes),  
and sample  $N(q)$  as many time as the number of unknown coefficients

# CUTS AND RESIDUES

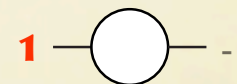
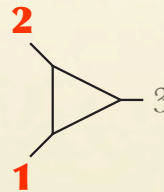
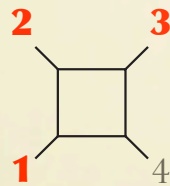
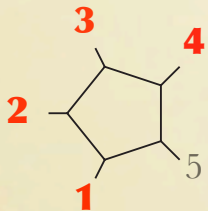
## cut-associated basis

For each cut  $(ijk\dots)$ ,  $D_i = D_j = D_k = \dots = 0$ , a basis of four massless vectors

$$\left\{ e_1^{(ijk\dots)}, e_2^{(ijk\dots)}, e_3^{(ijk\dots)}, e_4^{(ijk\dots)} \right\}$$

$$\begin{aligned} \left( e_i^{(ijk\dots)} \right)^2 &= 0, & e_1^{(ijk\dots)} \cdot e_3^{(ijk\dots)} &= e_1^{(ijk\dots)} \cdot e_4^{(ijk\dots)} = 0, \\ e_2^{(ijk\dots)} \cdot e_3^{(ijk\dots)} &= e_2^{(ijk\dots)} \cdot e_4^{(ijk\dots)} = 0, & e_1^{(ijk\dots)} \cdot e_2^{(ijk\dots)} &= -e_3^{(ijk\dots)} \cdot e_4^{(ijk\dots)} = 1 \end{aligned}$$

use independent external momenta + auxiliary orthogonal complement:



## 4-vectors vs components

- Loop momentum decomposition

$$q + p_i = \sum_{\alpha=1}^4 x_{\alpha} e_{\alpha}^{(ijk\dots)}$$

# THE SHAPE OF RESIDUES

legs	basis	
	external ( $p_i$ )	auxiliary ( $v_i$ )
5	4	0
4	3	1
3	2	2
2	1	3
1	0	4



## $\Delta$ -variables

- ISP's = Irreducible Scalar Products:
  - $q$ -components which can vary under cut-conditions
  - spurious: vanishing upon integration
  - non-spurious: non-vanishing upon integration  $\Rightarrow$  MI's
- @ 1-Loop
  - $(q \cdot p_i)$  are ALL reducible
  - ISP's could be chosen to be ALL spurious Pittau, de l'Aguila
  - $n$ -ple cut identifies an  $n$ -point diagram

# INTEGRAND-REDUCTION BEYOND ONE-LOOP

Ossola & P.M. (2011)

Badger, Frellesvig, Zhang (2011)

Zhang (2012)

Mirabella, Ossola, Peraro, & P.M (2012)

Kleiss, Malamos, Papadopoulos, Verheyne (2012)



# TWO-LOOP INTEGRAND REDUCTION

Ossola & P.M. (2011)

## Four-Dimensional Algorithm

$$\mathcal{A}_n = \int d^{4-2\epsilon} q \int d^{4-2\epsilon} k A(q, k), \quad A(q, k) = \frac{N(q, k)}{D_1 D_1 \cdots D_n}, \quad D_i = (\alpha_i q + \beta_i k + p_i)^2 - m_i^2, \quad \alpha_i, \beta_i \in \{0, 1\}$$

## educated guess: Master-Decomposition Formula (4-dim)

$$A(q, k) = \sum_{i_1 \ll i_8}^n \frac{\Delta_{i_1, \dots, i_8}(q, k)}{D_{i_1} D_{i_2} \cdots D_{i_8}} + \sum_{i_1 \ll i_7}^n \frac{\Delta_{i_1, \dots, i_7}(q, k)}{D_{i_1} D_{i_2} \cdots D_{i_7}} + \cdots + \sum_{i_1 \ll i_2}^n \frac{\Delta_{i_1, i_2}(q, k)}{D_{i_1} D_{i_2}}.$$

$$N(q, k) = \sum_{i_1 \ll i_8}^n \Delta_{i_1, \dots, i_8}(q, k) \prod_{h \neq i_1, \dots, i_8}^n D_h + \sum_{i_1 \ll i_7}^n \Delta_{i_1, \dots, i_7}(q, k) \prod_{h \neq i_1, \dots, i_7}^n D_h + \cdots + \sum_{i_1 \ll i_2}^n \Delta_{i_1, i_2}(q, k) \prod_{h \neq i_1, i_2}^n D_h, \quad (2.5)$$

In Dim-Reg higher-point higher-dim MI's can appear

□ Problem: what is the form of the residues?

 “find the right variables encoding the cut-structure”

# THE SHAPE OF RESIDUES

## **m-particle cut**

the vanishing of  $m$  denominators present in that diagram.

- Loop momentum decomposition

$$q + p_i = \sum_{\alpha=1}^4 x_{\alpha} e_{\alpha}^{(ijk\dots)},$$

$$k + p_i = \sum_{\alpha=1}^4 y_{\alpha} e_{\alpha}^{(ijk\dots)},$$

## **m-particle residue:** $\Delta_{i_1, \dots, i_m}$

legs	basis	
	external ( $p_i$ )	auxiliary ( $v_i$ )
5	4	0
4	3	1
3	2	2
2	1	3
1	0	4

## **ISP's variables in** $\Delta_{i_1, \dots, i_m}$

- ISP's = Irreducible Scalar Products:
  - spurious: vanishing upon integration
  - non-spurious: non-vanishing upon integration  $\Rightarrow$  MI's
- @ 2-Loop
  - $(q \cdot p_i)$  and  $(k \cdot p_i)$  can be ISP's ( $\neq$  1-Loop)
  - some ISP's could be chosen to be spurious
  - ISP's from:
    - \* direct inspection of the cut-solutions
    - \* relations among scalar products *via* Gram's Id'y

Badger, Frellesvig, Zhang

**MULTI-LOOP SCATTERING AMP'S  
FROM MULTIVARIATE POLYNOMIAL DIVISION**



# ALGEBRAIC GEOMETRY

- deals with multivariate polynomials in  $\mathbf{z} = (z_1, z_2, \dots)$ .
- **Ideal**  $\mathcal{J} \equiv \langle \omega_1(\mathbf{z}) \cdots \omega_s(\mathbf{z}) \rangle$  generated by  $\omega_i$ 
  - $\mathcal{J} = \{ \sum_i h_i(\mathbf{z}) \omega_i(\mathbf{z}) \}$
  - polynomial coefficients  $h_i(\mathbf{z})$
- **Multivariate polynomial division** of  $f(\mathbf{z})$  modulo  $\omega_1(\mathbf{z}), \dots, \omega_s(\mathbf{z})$ 
  - needs an order, i.e.  $z_1 z_2 \stackrel{?}{>} z_1^2$
  - $\rightsquigarrow f(\mathbf{z}) = \sum_i h_i(\mathbf{z}) \omega_i(\mathbf{z}) + \mathcal{R}(\mathbf{z})$
  - $h_i(\mathbf{z})$  &  $\mathcal{R}(\mathbf{z})$  not unique
- **Gröbner basis**  $\{g_1(\mathbf{z}), \dots, g_r(\mathbf{z})\}$ 
  - exists (Buchberger's algorithm) & generates  $\mathcal{J}$
  - $\rightsquigarrow$  unique  $\mathcal{R}(\mathbf{z})$
- **Hilbert's Nullstellensatz**
  - $V(\mathcal{J}) =$  set of common zeros of  $\mathcal{J}$
  - $(f = 0 \text{ in } V(\mathcal{J})) \Rightarrow (f^r \in \mathcal{J} \text{ for some } r)$
  - **Weak Nullstellensatz:**  $(V(\mathcal{J}) = \emptyset) \Leftrightarrow (1 \in \mathcal{J})$

# MULTIVARIATE POLYNOMIAL DIVISION

Zhang (2012);  
Mirabella, Ossola, Peraro, & P.M. (2012)

## Ideal

$$\mathcal{J}_{i_1 \dots i_n} = \langle D_{i_1}, \dots, D_{i_n} \rangle \equiv \left\{ \sum_{\kappa=1}^n h_{\kappa}(\mathbf{z}) D_{i_{\kappa}}(\mathbf{z}) : h_{\kappa}(\mathbf{z}) \in P[\mathbf{z}] \right\}$$

## Groebner Basis

$$\mathcal{G}_{i_1 \dots i_n} = \{g_1(\mathbf{z}), \dots, g_m(\mathbf{z})\}$$

$n$ -ple cut-conditions

$$D_{i_1} = \dots = D_{i_n} = 0 \quad \Leftrightarrow \quad g_1 = \dots = g_m = 0$$

# MULTIVARIATE POLYNOMIAL DIVISION

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$n$ -ple cut-conditions

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## Polynomial Division

$$\mathcal{N}_{i_1 \dots i_n}(\mathbf{z}) = \Gamma_{i_1 \dots i_n} + \Delta_{i_1 \dots i_n}(\mathbf{z}),$$

## Remainder ~ Residue

$$\Delta_{i_1 \dots i_n}(\mathbf{z})$$

## Quotients

$$\begin{aligned} \Gamma_{i_1 \dots i_n} &= \sum_{i=1}^m Q_i(\mathbf{z}) g_i(\mathbf{z}) && \text{belongs to the ideal } \mathcal{J}_{i_1 \dots i_n}, \\ &= \sum_{\kappa=1}^n \mathcal{N}_{i_1 \dots i_{\kappa-1} i_{\kappa+1} \dots i_n}(\mathbf{z}) D_{i_{\kappa}}(\mathbf{z}). \end{aligned}$$

# MULTI-LOOP RECURSIVE INTEGRAND REDUCTION

Mirabella, Ossola, Peraro, & P.M. (2012)

$$\mathcal{I}_{i_1 \dots i_n} = \sum_{\kappa=1}^k \mathcal{I}_{i_1 \dots i_{\kappa-1} i_{\kappa+1} i_n} + \frac{\Delta_{i_1 \dots i_n}}{D_{i_1} \dots D_{i_n}} .$$

n-denominator  
integrand

(n-1)-denominator  
integrand

remainder = residue



# REDUCIBILITY CRITERION

Mirabella, Ossola, Peraro, & P.M. (2012)

**Proposition 2.1.** *The integrand  $\mathcal{I}_{i_1 \dots i_n}$  is reducible iff the remainder of the division modulo a Gröbner basis vanishes, i.e. iff  $\mathcal{N}_{i_1 \dots i_n} \in \mathcal{J}_{i_1 \dots i_n}$ .*

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**Proposition 2.2.** *Any  $n$ -particle integrand with  $n > 4\ell$  is reducible.*

*Proof.* In this case, the system is over-constrained, namely the number  $n$  of equations is larger than the number  $4\ell$  of indeterminates. The  $n$  propagators cannot vanish simultaneously, i.e.

$$D_{i_1}(\mathbf{z}) = \dots = D_{i_n}(\mathbf{z}) = 0 \quad (2.7)$$

has no solution. Therefore, according to the weak Nullstellensatz theorem

$$1 = \sum_{\kappa=1}^n \omega_{\kappa}(\mathbf{z}) D_{i_{\kappa}}(\mathbf{z}) \in \mathcal{J}_{i_1 \dots i_n}, \quad (2.8)$$

for some  $\omega_{\kappa} \in P[\mathbf{z}]$ . Irrespective of the monomial order, a (reduced) Gröbner basis is  $\mathcal{G} = \{g_1\} = \{1\}$ . Eq. (2.5) becomes

$$\mathcal{N}_{i_1 \dots i_n}(\mathbf{z}) = \mathcal{N}_{i_1 \dots i_n}(\mathbf{z}) \times 1 \in \mathcal{J}_{i_1 \dots i_n}, \quad (2.9)$$

thus  $\mathcal{I}_{i_1 \dots i_n}$  is reducible. □

# ONE-LOOP INTEGRAND REDUCTION

In  $d$ -dimensions, the generic  $n$ -point one-loop integrand reads  $\mathcal{I}_{0\dots(n-1)} \equiv \frac{\mathcal{N}_{0\dots(n-1)}(q, \mu^2)}{D_0(q, \mu^2) \cdots D_{n-1}(q, \mu^2)}$ .

for each  $\mathcal{I}_{i_1\dots i_k}$  we define a basis  $\mathcal{E}^{(i_1\dots i_k)} = \{e_i\}_{i=1,\dots,4}$ .

If  $k \geq 5$ , then  $e_i = k_i$ , where  $k_i$  are four external momenta.

If  $k < 5$ , then  $e_i$  are chosen to fulfill the following relations:

$$\begin{aligned} e_1^2 = e_2^2 = 0, & & e_1 \cdot e_2 = 1, \\ e_3^2 = e_4^2 = \delta_{k4}, & & e_3 \cdot e_4 = -(1 - \delta_{k4}). \end{aligned}$$

In terms of  $\mathcal{E}^{(i_1\dots i_k)}$ , the loop momentum can be decomposed as,  $q^\mu = -p_{i_1}^\mu + x_1 e_1^\mu + x_2 e_2^\mu + x_3 e_3^\mu + x_4 e_4^\mu$ .

each numerator  $\mathcal{N}_{i_1\dots i_k}$  can be treated as a rank-  $k$  polynomial in  $\mathbf{z} \equiv (x_1, x_2, x_3, x_4, \mu^2)$ ,

$$\mathcal{N}_{i_1\dots i_k} = \sum_{\vec{j} \in J(k)} \alpha_{\vec{j}} z_1^{j_1} z_2^{j_2} z_3^{j_3} z_4^{j_4} z_5^{j_5},$$

$$J(k) \equiv \{\vec{j} = (j_1, \dots, j_5) : j_1 + j_2 + j_3 + j_4 + 2j_5 \leq k\}.$$

☑ *Step 1.* Since  $n > 5$ , the Proposition 2.2 guarantees that  $\mathcal{N}_{0\dots n-1}$  is reducible, and, by iteration, it can be written as a linear combination of 5-point integrands  $\mathcal{I}_{i_1\dots i_5}$ .

☑ *Step 1.* Since  $n > 5$ , the Proposition 2.2 guarantees that  $\mathcal{N}_{0\dots n-1}$  is reducible, and, by iteration, it can be written as a linear combination of 5-point integrands  $\mathcal{I}_{i_1\dots i_5}$ .

☑ *Step 2.* The numerator of each  $\mathcal{I}_{i_1\dots i_5}$  is a rank-5 polynomial in  $\mathbf{z}$ . We define the ideal  $\mathcal{J}_{i_1\dots i_5}$ , and compute the Gröbner basis  $\mathcal{G}_{i_1\dots i_5} = (g_1, \dots, g_5)$ , which is found to have a remarkably simple form:


$$g_i(\mathbf{z}) = c_i + z_i, \quad (i = 1, \dots, 5). \quad \text{[keep it in mind!]}$$

The division of  $\mathcal{N}_{i_1\dots i_5}$  modulo  $\mathcal{G}_{i_1\dots i_5}$  gives a *constant* remainder,

$$\Delta_{i_1\dots i_5} = c_0.$$

$$\Gamma_{i_1\dots i_5} = \sum_{\kappa=1}^5 \mathcal{N}_{i_1\dots i_{\kappa-1}i_{\kappa+1}\dots i_5}(\mathbf{z}) D_{i_\kappa}(\mathbf{z}),$$

where  $\mathcal{N}_{i_1\dots i_{\kappa-1}i_{\kappa+1}\dots i_5}$  are the numerators of the 4-point integrands,  $\mathcal{I}_{i_1\dots i_{\kappa-1}i_{\kappa+1}\dots i_5}$ , obtained by removing the  $i_\kappa$ -th denominator.

 *Step 3.* For each  $\mathcal{I}_{i_1 \dots i_4}$ , the numerator  $\mathcal{N}_{i_1 \dots i_4}$  is a rank-4 polynomial in  $\mathbf{z}$ . The Gröbner basis  $\mathcal{G}_{i_1 \dots i_4}$  of the ideal  $\mathcal{J}_{i_1 \dots i_4}$  contains four elements. Dividing  $\mathcal{N}_{i_1 \dots i_4}$  modulo  $\mathcal{G}_{i_1 \dots i_4}$ , we obtain

$$\Delta_{i_1 \dots i_4} = c_0 + c_1 x_4 + \mu^2 (c_2 + c_3 x_4 + \mu^2 c_4).$$

$$\Gamma_{i_1 \dots i_4} = \sum_{\kappa=1}^4 \mathcal{N}_{i_1 \dots i_{\kappa-1} i_{\kappa+1} \dots i_4}(\mathbf{z}) D_{i_{\kappa}}(\mathbf{z}),$$

contains the numerators of 3-point integrands  $\mathcal{I}_{i_1 \dots i_{\kappa-1} i_{\kappa+1} \dots i_4}$ .

☑ *Step 3.* For each  $\mathcal{I}_{i_1 \dots i_4}$ , the numerator  $\mathcal{N}_{i_1 \dots i_4}$  is a rank-4 polynomial in  $\mathbf{z}$ . The Gröbner basis  $\mathcal{G}_{i_1 \dots i_4}$  of the ideal  $\mathcal{J}_{i_1 \dots i_4}$  contains four elements. Dividing  $\mathcal{N}_{i_1 \dots i_4}$  modulo  $\mathcal{G}_{i_1 \dots i_4}$ , we obtain

$$\Delta_{i_1 \dots i_4} = c_0 + c_1 x_4 + \mu^2 (c_2 + c_3 x_4 + \mu^2 c_4) .$$

$$\Gamma_{i_1 \dots i_4} = \sum_{\kappa=1}^4 \mathcal{N}_{i_1 \dots i_{\kappa-1} i_{\kappa+1} \dots i_4}(\mathbf{z}) D_{i_{\kappa}}(\mathbf{z}) ,$$

contains the numerators of 3-point integrands  $\mathcal{I}_{i_1 \dots i_{\kappa-1} i_{\kappa+1} \dots i_4}$ .

☑ *Step 4.* The Gröbner basis  $\mathcal{G}_{i_1 i_2 i_3}$  is formed by three elements, and is used to divide  $\mathcal{N}_{i_1 i_2 i_3}$ . The remainder  $\Delta_{i_1 i_2 i_3}$  is polynomial in  $\mu^2$  and in the third and fourth components of  $q$  in the basis  $\mathcal{E}^{(i_1 i_2 i_3)}$ ,

$$\Delta_{i_1 i_2 i_3} = c_0 + c_1 x_3 + c_2 x_3^2 + c_3 x_3^3 + c_4 x_4 + c_5 x_4^2 + c_6 x_4^3 + \mu^2 (c_7 + c_8 x_3 + c_9 x_4) .$$

The term  $\Gamma_{i_1 i_2 i_3}$  generates the rank-2 numerators of the 2-point integrands  $\mathcal{I}_{i_1 i_2}$ ,  $\mathcal{I}_{i_1 i_3}$ , and  $\mathcal{I}_{i_2 i_3}$ .

☑ *Step 5.* The remainder of the division of  $\mathcal{N}_{i_1 i_2}$  by the two elements of  $\mathcal{G}_{i_1 i_2}$  is:

$$\Delta_{i_1 i_2} = c_0 + c_1 x_2 + c_2 x_3 + c_3 x_4 + c_4 x_2^2 + c_5 x_3^2 + c_6 x_4^2 + c_7 x_2 x_3 + c_8 x_2 x_4 + c_9 \mu^2 .$$

It is polynomial in  $\mu^2$  and in the last three components of  $q$  in the basis  $\mathcal{E}^{(i_1 i_2)}$ . The reducible term of the division,  $\Gamma_{i_1 i_2}$ , generates the rank-1 integrands,  $\mathcal{I}_{i_1}$ , and  $\mathcal{I}_{i_2}$ .



☑ *Step 5.* The remainder of the division of  $\mathcal{N}_{i_1 i_2}$  by the two elements of  $\mathcal{G}_{i_1 i_2}$  is:

$$\Delta_{i_1 i_2} = c_0 + c_1 x_2 + c_2 x_3 + c_3 x_4 + c_4 x_2^2 + c_5 x_3^2 + c_6 x_4^2 + c_7 x_2 x_3 + c_8 x_2 x_4 + c_9 \mu^2 .$$

It is polynomial in  $\mu^2$  and in the last three components of  $q$  in the basis  $\mathcal{E}^{(i_1 i_2)}$ . The reducible term of the division,  $\Gamma_{i_1 i_2}$ , generates the rank-1 integrands,  $\mathcal{I}_{i_1}$ , and  $\mathcal{I}_{i_2}$ .

☑ *Step 6.* The numerator of the 1-point integrands is linear in the components of the loop momentum in the basis  $\mathcal{E}^{(i_1)}$ ,

$$\mathcal{N}_{i_1} = \beta_0 + \sum_{j=1}^4 \beta_j x_j .$$

The only element of the Gröbner basis  $\mathcal{G}_{i_1}$  is  $D_{i_1}$ , which is quadratic in  $\mathbf{z}$ . Therefore the division modulo  $\mathcal{G}_{i_1}$ , leads to a vanishing quotient, hence

$$\mathcal{N}_{i_1} = \Delta_{i_1} .$$

☑ *Step 7.* Collecting all the remainders computed in the previous steps, we obtain the complete decomposition of  $\mathcal{I}_{0\dots n-1}$  in terms of its multi-pole structure

$$\mathcal{I}_{0\dots n-1} = \sum_{k=1}^5 \left( \sum_{1=i_1 < \dots < i_k}^{n-1} \frac{\Delta_{i_1 \dots i_k}}{D_{i_1} \cdots D_{i_k}} \right) .$$

which reproduces the well-known one-loop  $d$ -dimensional integrand decomposition formula

Ossola, Papadopoulos, Pittau  
Ellis, Giele, Kunszt, Melnikov

----- SPINORS @ MATHEMATICA (S@M) -----

Version: S@M 1.0 (3-APR-2007)

Authors:

Daniel Maitre (SLAC),  
Pierpaolo Mastrolia (University of Zurich)

A list of all functions provided by the package  
is stored in the variable  
\$SpinorsFunctions

**DeclareSpinor**[p[1], p[2], p[3], p[4], p[5]]

{p[1], p[2], p[3], p[4], p[5]} added to the list of spinors

**GenMomenta**[[p[1], p[2], p[3], p[4], p[5]]]

Momenta for the spinors p[1], p[2], p[3], p[4], p[5] generated.

**DeclareLVector**[q, qtemp]

{q, qtemp} added to the list of Lorentz vectors

**DeclareLVector**[e[1], e[2], e[3], e[4]]

{e[1], e[2], e[3], e[4]} added to the list of Lorentz vectors

---

**GroebnerBasis**[[*poly*<sub>1</sub>, *poly*<sub>2</sub>, ...], {*x*<sub>1</sub>, *x*<sub>2</sub>, ...}] gives a  
list of polynomials that form a Gröbner basis for the set of polynomials *poly*<sub>*i*</sub>.

---

**PolynomialReduce**[*poly*, {*poly*<sub>1</sub>, *poly*<sub>2</sub>, ...}, {*x*<sub>1</sub>, *x*<sub>2</sub>, ...}] yields  
a list representing a reduction of *poly* in terms of the *poly*<sub>*i*</sub>. The list has the form  
{*a*<sub>1</sub>, *a*<sub>2</sub>, ...}, *b*}, where *b* is minimal and *a*<sub>1</sub> *poly*<sub>1</sub> + *a*<sub>2</sub> *poly*<sub>2</sub> + ... + *b* is exactly *poly*. >>

## 5ple-cut

### ■ q-decomposition

```
qdeco = Sum[x[i] * p[i], {i, 1, 4}]
```

```
p[1] x[1] + p[2] x[2] + p[3] x[3] + p[4] x[4]
```

### ■ Explicit expression of denominators in terms of x[i]

```
For[i = 0, i ≤ 4, i++,  
  Dnr[i] = Den[i];  
  Print["D_", i, " = ", Dnr[i]];  
];
```

```
D_0 = -μ2 + MP[q, q]
```

```
D_1 = -μ2 + MP[q, q] + 2 MP[q, p[1]]
```

```
D_2 = -μ2 + MP[q, q] + 2 MP[q, p[1]] + 2 MP[q, p[2]] + 2 MP[p[1], p[2]]
```

```
D_3 = -μ2 + MP[q, q] + 2 MP[q, p[1]] + 2 MP[q, p[2]] +  
  2 MP[q, p[3]] + 2 MP[p[1], p[2]] + 2 MP[p[1], p[3]] + 2 MP[p[2], p[3]]
```

```
D_4 = -μ2 + MP[q, q] - 2 MP[q, p[5]]
```

```
D_0 = -1. μ2 + 3.89861 x[1.] x[2.] + 10.8202 x[1.] x[3.] +  
  6.38478 x[2.] x[3.] - 2.95514 x[1.] x[4.] - 3.47983 x[2.] x[4.] - 14.6687 x[3.] x[4.]
```

```
D_1 = -1. μ2 + 3.89861 x[2.] + 3.89861 x[1.] x[2.] +  
  10.8202 x[3.] + 10.8202 x[1.] x[3.] + 6.38478 x[2.] x[3.] - 2.95514 x[4.] -  
  2.95514 x[1.] x[4.] - 3.47983 x[2.] x[4.] - 14.6687 x[3.] x[4.]
```

```
D_2 = 3.89861 - 1. μ2 + 3.89861 x[1.] + 3.89861 x[2.] + 3.89861 x[1.] x[2.] +  
  17.205 x[3.] + 10.8202 x[1.] x[3.] + 6.38478 x[2.] x[3.] - 6.43497 x[4.] -  
  2.95514 x[1.] x[4.] - 3.47983 x[2.] x[4.] - 14.6687 x[3.] x[4.]
```

```
D_3 = 21.1036 - 1. μ2 + 14.7189 x[1.] + 10.2834 x[2.] + 3.89861 x[1.] x[2.] +  
  17.205 x[3.] + 10.8202 x[1.] x[3.] + 6.38478 x[2.] x[3.] - 21.1036 x[4.] -  
  2.95514 x[1.] x[4.] - 3.47983 x[2.] x[4.] - 14.6687 x[3.] x[4.]
```

```
D_4 = -1. μ2 + 11.7637 x[1.] + 6.80356 x[2.] + 3.89861 x[1.] x[2.] +  
  2.53636 x[3.] + 10.8202 x[1.] x[3.] + 6.38478 x[2.] x[3.] - 21.1036 x[4.] -  
  2.95514 x[1.] x[4.] - 3.47983 x[2.] x[4.] - 14.6687 x[3.] x[4.]
```

```

Vars = Union[Table[x[i], {i, 1, 4}]] // N;
Vars = Append[Vars,  $\mu^2$ ]
{x[1.], x[2.], x[3.], x[4.],  $\mu^2$ }

```

```

GB = GroebnerBasis[Cut[0, 1, 2, 3, 4], Vars];

```

```

For[i = 1, i ≤ Length[GB], i++,
  Print[i, " : ", GB[[i]]];
]

```

```

1 : 31.9034 + 1.  $\mu^2$ 

```

```

2 : 3.44658 + 1. x[4.]

```

```

3 : -0.196787 + 1. x[3.]

```

```

4 : 3.15867 + 1. x[2.]

```

```

5 : 4.39864 + 1. x[1.]

```

```

Num5 = (
  c[0] +
  Sum[c[i1] * Vars[[i1]], {i1, 1, Length[Vars]}] +
  Sum[c[i1, i2] * Vars[[i1]] * Vars[[i2]],
    {i1, 1, Length[Vars]}, {i2, i1, Length[Vars]}] +
  Sum[c[i1, i2, i3] * Vars[[i1]] * Vars[[i2]] * Vars[[i3]],
    {i1, 1, Length[Vars]}, {i2, i1, Length[Vars]}, {i3, i2, Length[Vars]}] +
  Sum[c[i1, i2, i3, i4] * Vars[[i1]] * Vars[[i2]] * Vars[[i3]] * Vars[[i4]],
    {i1, 1, Length[Vars]}, {i2, i1, Length[Vars]},
    {i3, i2, Length[Vars]}, {i4, i3, Length[Vars]}] +
  Sum[c[i1, i2, i3, i4, i5] * Vars[[i1]] * Vars[[i2]] * Vars[[i3]] *
    Vars[[i4]] * Vars[[i5]], {i1, 1, Length[Vars]}, {i2, i1, Length[Vars]},
    {i3, i2, Length[Vars]}, {i4, i3, Length[Vars]}, {i5, i4, Length[Vars]}]
)

```

```

c[0] +  $\mu^2$  c[5] +  $\mu^2^2$  c[5, 5] +  $\mu^2^3$  c[5, 5, 5] +  $\mu^2^4$  c[5, 5, 5, 5] +  $\mu^2^5$  c[5, 5, 5, 5, 5] +
c[1] x[1.] +  $\mu^2$  c[1, 5] x[1.] +  $\mu^2^2$  c[1, 5, 5] x[1.] +  $\mu^2^3$  c[1, 5, 5, 5] x[1.] +
 $\mu^2^4$  c[1, 5, 5, 5, 5] x[1.] + c[1, 1] x[1.]2 +  $\mu^2$  c[1, 1, 5] x[1.]2 +  $\mu^2^2$  c[1, 1, 5, 5] x[1.]2 +
 $\mu^2^3$  c[1, 1, 5, 5, 5] x[1.]2 + c[1, 1, 1] x[1.]3 +  $\mu^2$  c[1, 1, 1, 5] x[1.]3 +
 $\mu^2^2$  c[1, 1, 1, 5, 5] x[1.]3 + c[1, 1, 1, 1] x[1.]4 +  $\mu^2$  c[1, 1, 1, 1, 5] x[1.]4 +
c[1, 1, 1, 1, 1] x[1.]5 + c[2] x[2.] +  $\mu^2$  c[2, 5] x[2.] +  $\mu^2^2$  c[2, 5, 5] x[2.] +
 $\mu^2^3$  c[2, 5, 5, 5] x[2.] +  $\mu^2^4$  c[2, 5, 5, 5, 5] x[2.] + c[1, 2] x[1.] x[2.] + ...

```

```

Length[%]

```

■ **5-point residue**

```
resto = PolynomialReduce[Num5, GB, Vars][[2]];
```

```
Length[MonomialList[resto, Vars]]
```

1

## 4ple-cut

### ■ q-decomposition

`qdeco = Sum[x[i] * e[i], {i, 1, 2}] + x[3] * v + x[4] * vperp`

$$e[1] x[1] + e[2] x[2] + \frac{1}{2} (e[3] + e[4]) x[3] + \frac{1}{2} (e[3] - e[4]) x[4]$$

Note that  $(e[3]-e[4])/2 = vperp$ , because it is orthogonal to  $e[1], e[2]$  and  $v=(e[3]+e[4])/2$

### ■ Explicit expression of denominators in terms of x[i]

$$D_0 = -1. \mu^2 - 3.92177 x[1.] x[2.] + 0.980442 x[3.]^2 - 0.980442 x[4.]^2$$

$$D_1 = -1. \mu^2 - 3.92177 x[2.] - 3.92177 x[1.] x[2.] + 0.980442 x[3.]^2 - 0.980442 x[4.]^2$$

$$D_2 = 12.4329 - 1. \mu^2 + 12.4329 x[1.] - 7.85602 x[2.] - 3.92177 x[1.] x[2.] + (0. + 4.26815 i) x[3.] + 0.980442 x[3.]^2 + 5.54049 x[4.] - 0.980442 x[4.]^2$$

$$D_3 = 4.43686 - 1. \mu^2 + 11.6475 x[1.] - 4.43686 x[2.] - 3.92177 x[1.] x[2.] + (0. + 2.66221 i) x[3.] + 0.980442 x[3.]^2 + 5.21447 x[4.] - 0.980442 x[4.]^2$$

$$D_4 = -1. \mu^2 + 3.92177 x[1.] - 3.92177 x[1.] x[2.] + 0.980442 x[3.]^2 - 0.980442 x[4.]^2$$

`GB = GroebnerBasis[Cut[0, 1, 2, 3], Vars];`

1 :  $29.1065 + 1.01962 \mu^2 - 0.193505 x[4.] + 1. x[4.]^2$

2 :  $(0. - 5.39591 i) + 1. x[3.] + (0. + 0.0179364 i) x[4.]$

3 :  $1. x[2.]$

4 :  $-0.852388 + 1. x[1.] + 0.451789 x[4.]$

### Num4

$$\begin{aligned} & c[0] + \mu^2 c[5] + \mu^2 c[5, 5] + c[1] x[1.] + \mu^2 c[1, 5] x[1.] + c[1, 1] x[1.]^2 + \\ & \mu^2 c[1, 1, 5] x[1.]^2 + c[1, 1, 1] x[1.]^3 + c[1, 1, 1, 1] x[1.]^4 + c[2] x[2.] + \\ & \mu^2 c[2, 5] x[2.] + c[1, 2] x[1.] x[2.] + \mu^2 c[1, 2, 5] x[1.] x[2.] + c[1, 1, 2] x[1.]^2 x[2.] + \\ & c[1, 1, 1, 2] x[1.]^3 x[2.] + c[2, 2] x[2.]^2 + \mu^2 c[2, 2, 5] x[2.]^2 + c[1, 2, 2] x[1.] x[2.]^2 + \\ & c[1, 1, 2, 2] x[1.]^2 x[2.]^2 + c[2, 2, 2] x[2.]^3 + c[1, 2, 2, 2] x[1.] x[2.]^3 + \\ & c[2, 2, 2, 2] x[2.]^4 + c[3] x[3.] + \mu^2 c[3, 5] x[3.] + c[1, 3] x[1.] x[3.] + \\ & \mu^2 c[1, 3, 5] x[1.] x[3.] + c[1, 1, 3] x[1.]^2 x[3.] + c[1, 1, 1, 3] x[1.]^3 x[3.] + \dots \end{aligned}$$

`Length[%]`

■ 4-point residue

```
resto = PolynomialReduce[Num4, GB, Vars, MonomialOrder → Lexicographic][[2]];
Length[MonomialList[resto, Vars]]
```

5



## 3ple-cut

### ■ q-decomposition

```
qdeco = Sum[x[i] * e[i], {i, 1, 4}]
```

```
e[1] x[1] + e[2] x[2] + e[3] x[3] + e[4] x[4]
```

### ■ Explicit expression of denominators in terms of x[i]

```
D_0 = -1. μ2 - 11.7637 x[1.] x[2.] + 11.7637 x[3.] x[4.]
```

```
D_1 = -1. μ2 - 11.7637 x[2.] - 11.7637 x[1.] x[2.] + 11.7637 x[3.] x[4.]
```

```
D_2 = 3.89861 - 1. μ2 + 3.89861 x[1.] - 18.5673 x[2.] - 11.7637 x[1.] x[2.] -  
(0.432545 - 5.13199 i) x[3.] + (0.432545 + 5.13199 i) x[4.] + 11.7637 x[3.] x[4.]
```

```
D_3 = 21.1036 - 1. μ2 + 14.7189 x[1.] - 21.1036 x[2.] - 11.7637 x[1.] x[2.] +  
(4.65563 + 6.37882 i) x[3.] - (4.65563 - 6.37882 i) x[4.] + 11.7637 x[3.] x[4.]
```

```
D_4 = -1. μ2 + 11.7637 x[1.] - 11.7637 x[1.] x[2.] + 11.7637 x[3.] x[4.]
```

```
GB = GroebnerBasis[Cut[0, 1, 4], Vars];
```

```
For[i = 1, i ≤ Length[GB], i++,  
  Print[i, " : ", GB[[i]]];  
]
```

```
1 : -0.0850072 μ2 + 1. x[3.] x[4.]
```

```
2 : 1. x[2.]
```

```
3 : 1. x[1.]
```

**Num3**

```
c[0] + μ2 c[5] + c[1] x[1.] + μ2 c[1, 5] x[1.] + c[1, 1] x[1.]2 + c[1, 1, 1] x[1.]3 +  
c[2] x[2.] + μ2 c[2, 5] x[2.] + c[1, 2] x[1.] x[2.] + c[1, 1, 2] x[1.]2 x[2.] +  
c[2, 2] x[2.]2 + c[1, 2, 2] x[1.] x[2.]2 + c[2, 2, 2] x[2.]3 + c[3] x[3.] +  
μ2 c[3, 5] x[3.] + c[1, 3] x[1.] x[3.] + c[1, 1, 3] x[1.]2 x[3.] + c[2, 3] x[2.] x[3.] +  
c[1, 2, 3] x[1.] x[2.] x[3.] + c[2, 2, 3] x[2.]2 x[3.] + c[3, 3] x[3.]2 +  
c[1, 3, 3] x[1.] x[3.]2 + c[2, 3, 3] x[2.] x[3.]2 + c[3, 3, 3] x[3.]3 +  
c[4] x[4.] + μ2 c[4, 5] x[4.] + c[1, 4] x[1.] x[4.] + c[1, 1, 4] x[1.]2 x[4.] +  
c[2, 4] x[2.] x[4.] + c[1, 2, 4] x[1.] x[2.] x[4.] + c[2, 2, 4] x[2.]2 x[4.] +  
c[3, 4] x[3.] x[4.] + c[1, 3, 4] x[1.] x[3.] x[4.] + c[2, 3, 4] x[2.] x[3.] x[4.] +  
c[3, 3, 4] x[3.]2 x[4.] + c[4, 4] x[4.]2 + c[1, 4, 4] x[1.] x[4.]2 +  
c[2, 4, 4] x[2.] x[4.]2 + c[3, 4, 4] x[3.] x[4.]2 + c[4, 4, 4] x[4.]3
```

```
Length[%]
```

### ■ 3-point residue

```
resto = PolynomialReduce[Num3, GB, Vars, MonomialOrder → Lexicographic][[2]];
```

```
Length[MonomialList[resto, Vars]]
```

```
10
```

#### **Nresto**

```
c[0.] + 1. μ2 c[5.] + 0.0850072 μ2 c[3., 4.] + c[3.] x[3.] +  
1. μ2 c[3., 5.] x[3.] + 0.0850072 μ2 c[3., 3., 4.] x[3.] + c[3., 3.] x[3.]2 +  
c[3., 3., 3.] x[3.]3 + c[4.] x[4.] + 1. μ2 c[4., 5.] x[4.] +  
0.0850072 μ2 c[3., 4., 4.] x[4.] + c[4., 4.] x[4.]2 + c[4., 4., 4.] x[4.]3
```

## 2ple-cut

### ■ q-decomposition

```
qdeco = Sum[x[i] * e[i], {i, 1, 4}]
```

```
e[1] x[1] + e[2] x[2] + e[3] x[3] + e[4] x[4]
```

### ■ Explicit expression of denominators in terms of x[i]

```
D_0 = -1. μ2 - 11.7637 x[1.] x[2.] + 11.7637 x[3.] x[4.]
```

```
D_1 = -1. μ2 - 11.7637 x[2.] - 11.7637 x[1.] x[2.] + 11.7637 x[3.] x[4.]
```

```
D_2 = 3.89861 - 1. μ2 + 3.89861 x[1.] - 18.5673 x[2.] - 11.7637 x[1.] x[2.] -  
(0.432545 - 5.13199 i) x[3.] + (0.432545 + 5.13199 i) x[4.] + 11.7637 x[3.] x[4.]
```

```
D_3 = 21.1036 - 1. μ2 + 14.7189 x[1.] - 21.1036 x[2.] - 11.7637 x[1.] x[2.] +  
(4.65563 + 6.37882 i) x[3.] - (4.65563 - 6.37882 i) x[4.] + 11.7637 x[3.] x[4.]
```

```
D_4 = -1. μ2 + 11.7637 x[1.] - 11.7637 x[1.] x[2.] + 11.7637 x[3.] x[4.]
```

```
GB = GroebnerBasis[Cut[0, 4], Vars];
```

```
1 : -0.0850072 μ2 + 1. x[3.] x[4.]
```

```
2 : 1. x[1.]
```

**Num2**

```
c[0] + μ2 c[5] + c[1] x[1.] + c[1, 1] x[1.]2 + c[2] x[2.] + c[1, 2] x[1.] x[2.] + c[2, 2] x[2.]2 +  
c[3] x[3.] + c[1, 3] x[1.] x[3.] + c[2, 3] x[2.] x[3.] + c[3, 3] x[3.]2 + c[4] x[4.] +  
c[1, 4] x[1.] x[4.] + c[2, 4] x[2.] x[4.] + c[3, 4] x[3.] x[4.] + c[4, 4] x[4.]2
```

```
Length[MonomialList[Num2, Vars]]
```

```
16
```

### ■ 2-point residue

```
resto = PolynomialReduce[Num2, GB, Vars, MonomialOrder → Lexicographic][[2]];
```

```
Length[MonomialList[resto, Vars]]
```

```
10
```

**Nresto**

```
c[0.] + 1. μ2 c[5.] + 0.0850072 μ2 c[3., 4.] +  
c[2.] x[2.] + c[2., 2.] x[2.]2 + c[3.] x[3.] + c[2., 3.] x[2.] x[3.] +  
c[3., 3.] x[3.]2 + c[4.] x[4.] + c[2., 4.] x[2.] x[4.] + c[4., 4.] x[4.]2
```

# 1ple-cut

## ■ q-decomposition

```
qdeco = Sum[x[i] * e[i], {i, 1, 4}]
```

```
e[1] x[1] + e[2] x[2] + e[3] x[3] + e[4] x[4]
```

## ■ Explicit expression of denominators in terms of x[i]

```
D_0 = -1. μ2 - 11.7637 x[1.] x[2.] + 11.7637 x[3.] x[4.]
```

```
D_1 = -1. μ2 - 11.7637 x[2.] - 11.7637 x[1.] x[2.] + 11.7637 x[3.] x[4.]
```

```
D_2 = 3.89861 - 1. μ2 + 3.89861 x[1.] - 18.5673 x[2.] - 11.7637 x[1.] x[2.] -  
(0.432545 - 5.13199 i) x[3.] + (0.432545 + 5.13199 i) x[4.] + 11.7637 x[3.] x[4.]
```

```
D_3 = 21.1036 - 1. μ2 + 14.7189 x[1.] - 21.1036 x[2.] - 11.7637 x[1.] x[2.] +  
(4.65563 + 6.37882 i) x[3.] - (4.65563 - 6.37882 i) x[4.] + 11.7637 x[3.] x[4.]
```

```
D_4 = -1. μ2 + 11.7637 x[1.] - 11.7637 x[1.] x[2.] + 11.7637 x[3.] x[4.]
```

```
GB = GroebnerBasis[Cut[0], Vars];
```

```
For[i = 1, i ≤ Length[GB], i++,  
  Print[i, " : ", GB[[i]]];  
]
```

```
1 : 0.0850072 μ2 + 1. x[1.] x[2.] - 1. x[3.] x[4.]
```

```
Num1
```

```
c[0] + c[1] x[1.] + c[2] x[2.] + c[3] x[3.] + c[4] x[4.]
```

```
Length[MonomialList[Num1, Vars]]
```

```
5
```

## ■ 1-point residue

```
resto = PolynomialReduce[Num1, GB, Vars, MonomialOrder → Lexicographic][[2]];
```

```
Length[MonomialList[resto, Vars]]
```

```
5
```

```
Nresto
```

```
c[0.] + c[1.] x[1.] + c[2.] x[2.] + c[3.] x[3.] + c[4.] x[4.]
```

```
quoz = PolynomialReduce[Num1, GB, Vars, MonomialOrder → Lexicographic][[1]]
```

```
{0.}
```



□ What can we do within this new framework?

# THE MAXIMUM-CUT THEOREM

Mirabella, Ossola, Peraro, & P.M. (2012)

At  $\ell$  loops, in four dimensions, we define a *maximum-cut* as a  $(4\ell)$ -ple cut

$$D_{i_1} = D_{i_2} = \cdots = D_{i_{4\ell}} = 0 ,$$

which constrains completely the components of the loop momenta. In four dimensions this implies the presence of four constraints for each loop momenta.

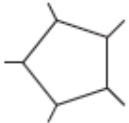
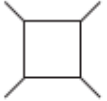
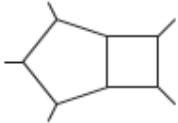
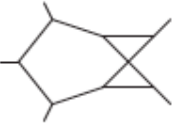

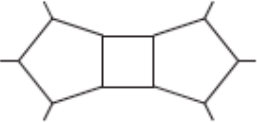
We assume that:

in non-exceptional phase-space points, a maximum-cut has a finite number  $n_s$  of solutions, each with multiplicity one.

Under this assumption we have the following

**Theorem 4.1** (Maximum cut). *The residue at the maximum-cut is a polynomial parametrised by  $n_s$  coefficients, which admits a univariate representation of degree  $(n_s - 1)$ .*

# EXAMPLES OF MAXIMUM-CUTS

diagram	$\Delta$	$n_s$	diagram	$\Delta$	$n_s$
	$c_0$	1		$c_0 + c_1 z$	2
	$\sum_{i=0}^3 c_i z^i$	4		$\sum_{i=0}^3 c_i z^i$	4
	$\sum_{i=0}^7 c_i z^i$	8		$\sum_{i=0}^7 c_i z^i$	8

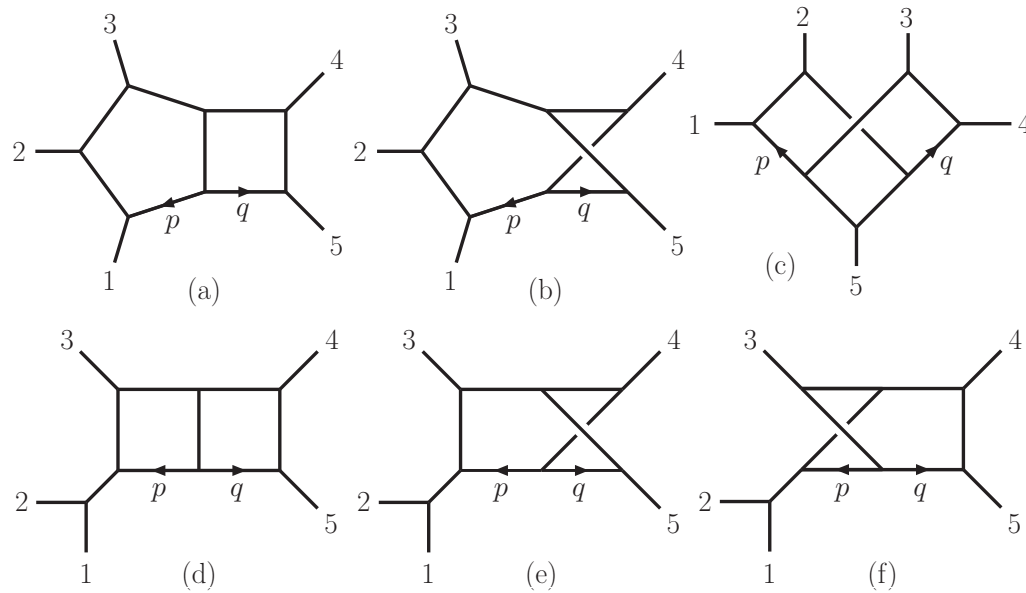
# 2-LOOP 5-POINT AMPLITUDES IN $N=4$ SYM

Bern, Czakon, Kosower, Roiban, Smirnov

Arkani-Hamed, Bourjaily, Cachazo, Caron-Houot, Trnka

Drummond, Henn, Trnka

Carrasco, Johansson





## Integrand

$$\mathcal{I}_{1\dots 8} \equiv \frac{\mathcal{N}_{1\dots 8}(q, k)}{D_1(q, k) \cdots D_8(q, k)},$$

## Momentum basis

$$q^\mu = \sum_{i=1}^4 y_i \tau_i^\mu, \quad k^\mu = \sum_{i=1}^4 x_i e_i^\mu.$$

## Generic Numerator

$$\mathcal{N}(x_1, x_2, x_3, x_4, y_1, y_2, y_3, y_4) = \sum_{\vec{j} \in J(k)} \alpha_{\vec{j}} x_1^{j_1} x_2^{j_2} x_3^{j_3} x_4^{j_4} y_1^{j_5} y_2^{j_6} y_3^{j_7} y_4^{j_8},$$

with  $J(k)$  being the set of values for the exponents compatible with the renormalizability

## Polynomial Division

$$\mathcal{N}_{i_1 \dots i_n}(\mathbf{z}) = \sum_{\kappa=1}^n \mathcal{N}_{i_1 \dots i_{\kappa-1} i_{\kappa+1} \dots i_n}(\mathbf{z}) D_{i_\kappa}(\mathbf{z}) + \Delta_{i_1 \dots i_n}(\mathbf{z}).$$

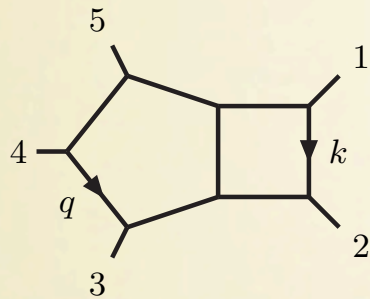
## 2-Loop Integrand Decomposition Formula (4D)

$$\mathcal{I}_n = \sum_{i_1 << i_8=1}^n \frac{\Delta_{i_1 \dots i_8}}{D_{i_1} \cdots D_{i_8}} + \sum_{i_1 << i_7=1}^n \frac{\Delta_{i_1 \dots i_7}}{D_{i_1} \cdots D_{i_7}} + \cdots + \sum_{i_1 < i_2=1}^n \frac{\Delta_{i_1 i_2}}{D_{i_1} D_{i_2}} + \sum_{i=1}^n \frac{\Delta_i}{D_i} + \mathcal{Q}_\emptyset$$

# THE PENTABOX DIAGRAM IN N=4 SYM

Ossola & P.M. (2011)

Mirabella, Ossola, Peraro, & P.M. (in progress)

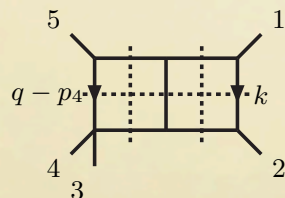
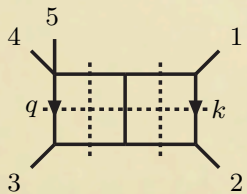
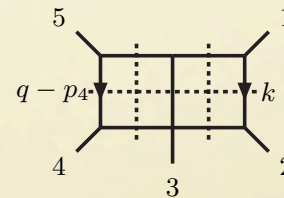
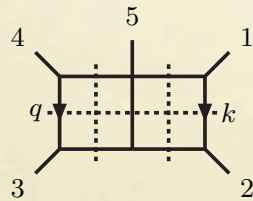
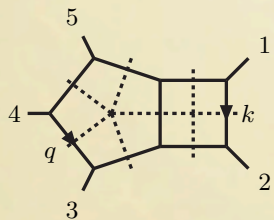


$$\begin{aligned}
 D_1 &= k^2 \\
 D_2 &= (k + p_2)^2 \\
 D_3 &= (k - p_1)^2 \\
 D_4 &= q^2 \\
 D_5 &= (q + p_3)^2 \\
 D_6 &= (q - p_4)^2 \\
 D_7 &= (q - p_4 - p_5)^2 \\
 D_8 &= (q + k + p_2 + p_3)^2 .
 \end{aligned}$$

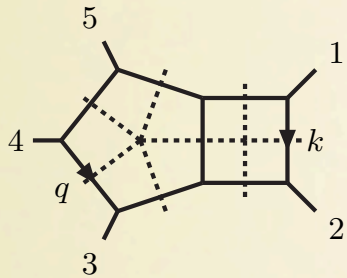
$$N(q, k) = 2q \cdot v + \alpha \quad \text{Carrasco & Johansson (2011)}$$

$$v^\mu = \frac{1}{4} \left( \gamma_{12}(p_1^\mu - p_2^\mu) + \gamma_{23}(p_2^\mu - p_3^\mu) + 2\gamma_{45}(p_4^\mu - p_5^\mu) + \gamma_{13}(p_1^\mu - p_3^\mu) \right)$$

$$\alpha = \frac{1}{4} \left( 2\gamma_{12}(s_{45} - s_{12}) + \gamma_{23}(s_{45} + 3s_{12} - s_{13}) + 2\gamma_{45}(s_{14} - s_{15}) + \gamma_{13}(s_{12} + s_{45} - s_{13}) \right)$$



# 5-POINT 8FOLD-CUT $D_1 = \dots = D_8 = 0$



$$\Delta_{12345678}(q, k) = \text{Res}_{12345678} \left\{ \mathcal{N}_{1\dots 8}(q, k) \right\} .$$

$$q^\mu = \sum_{i=1}^4 y_i \tau_i^\mu , \quad k^\mu = \sum_{i=1}^4 x_i e_i^\mu . \quad e_1 = p_1, \quad e_2 = p_2, \quad \tau_1 = p_3, \quad \tau_2 = p_4 .$$

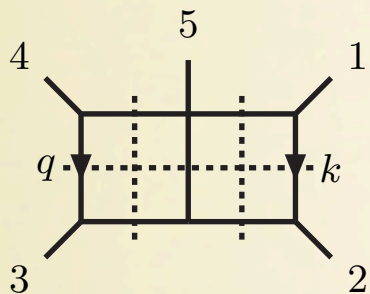
$$\Delta_{12345678}(q, k) = c_{12345678,0} + c_{12345678,1} y_4 + c_{12345678,2} x_3 + c_{12345678,3} x_4 .$$

[Maximum Cut Thm]

generic residue

# 5-POINT 7FOLD-CUT

$$D_1 = \dots = D_6 = D_8 = 0$$



$$\Delta_{1234568}(q, k) = \text{Res}_{1234568} \left\{ \frac{N(q, k) - \Delta_{12345678}(q, k)}{D_7} \right\}.$$

$$q^\mu = \sum_{i=1}^4 y_i \tau_i^\mu, \quad k^\mu = \sum_{i=1}^4 x_i e_i^\mu.$$

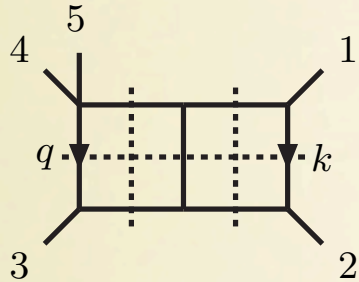
$$e_1 = p_1, \quad e_2 = p_2, \quad \tau_1 = p_3, \quad \tau_2 = p_4.$$

$$\begin{aligned} \Delta_{1234568} = & c_0 + c_1 x_3 + c_2 x_3^2 + c_3 x_3^3 + c_4 x_3^4 + c_5 x_4 + c_6 x_4^2 + c_7 x_4^3 + c_8 x_4^4 \\ & + c_9 y_3 + c_{10} x_4 y_3 + c_{11} y_3^2 + c_{12} x_4 y_3^2 + c_{13} y_3^3 + c_{14} x_4 y_3^3 + c_{15} y_3^4 \\ & + c_{16} x_4 y_3^4 + c_{17} y_4 + c_{18} x_3 y_4 + c_{19} x_3^2 y_4 + c_{20} x_3^3 y_4 + c_{21} x_3^4 y_4 + c_{22} x_4 y_4 \\ & + c_{23} x_4^2 y_4 + c_{24} x_4^3 y_4 + c_{25} x_4^4 y_4 + c_{26} y_4^2 + c_{27} x_4 y_4^2 + c_{28} y_4^3 + c_{29} x_4 y_4^3 \\ & + c_{30} y_4^4 + c_{31} x_4 y_4^4. \end{aligned} \tag{3.18}$$

generic residue

# 4-POINT 7FOLD-CUT

$$D_1 = \dots = D_5 = D_7 = D_8 = 0.$$



$$\Delta_{1234578}(q, k) = \text{Res}_{1234578} \left\{ \frac{N(q, k) - \Delta_{12345678}(q, k)}{D_6} \right\},$$

$$\begin{aligned} e_1^\mu &= p_1^\mu, & e_2^\mu &= p_2^\mu, \\ \tau_1^\mu &= p_3^\mu, & \tau_2^\mu &= P_{45}^\mu - \frac{s_{45}}{2P_{45} \cdot \tau_1} \tau_1^\mu. \end{aligned}$$

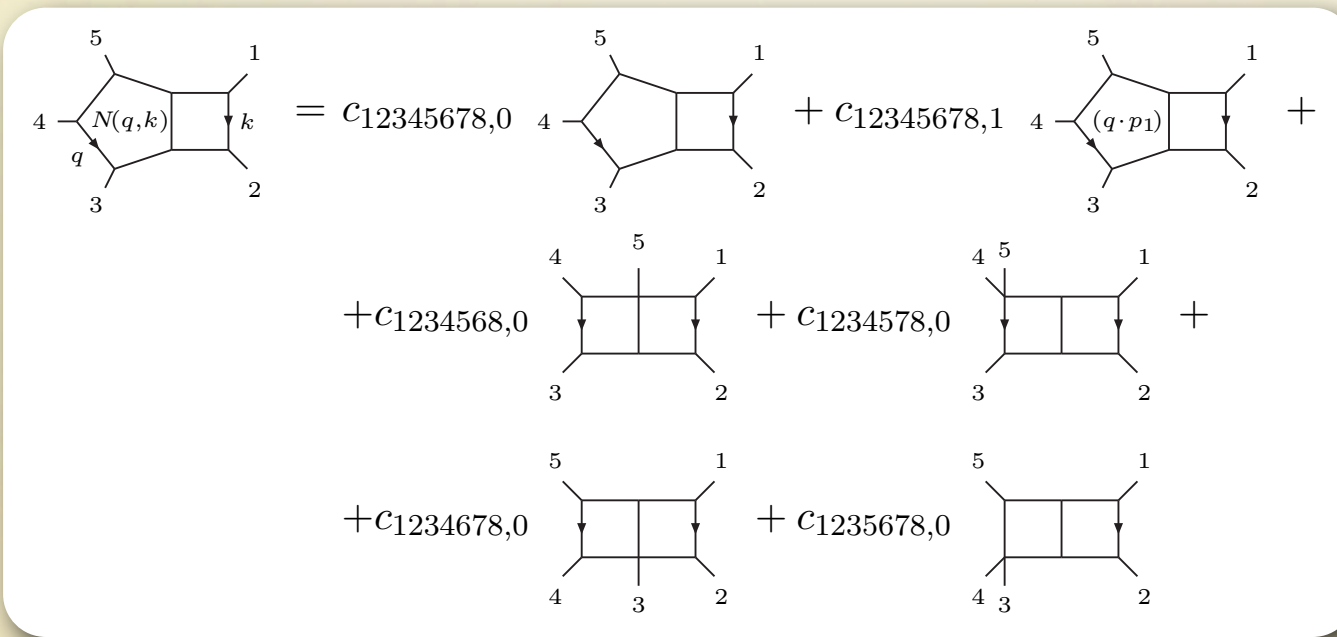
parametrized using thirty-two monomials

$$\{1, x_3, x_3^2, x_3^3, x_3^4, x_4, x_4^2, x_4^3, x_4^4, y_3, x_4 y_3, y_3^2, x_4 y_3^2, y_3^3, x_4 y_3^3, y_3^4, x_4 y_3^4, y_4, x_3 y_4, x_3^2 y_4, x_3^3 y_4, x_3^4 y_4, x_4 y_4, x_4^2 y_4, x_4^3 y_4, x_4^4 y_4, y_4^2, x_4 y_4^2, y_4^3, x_4 y_4^3, y_4^4, x_4 y_4^4\}.$$

generic residue

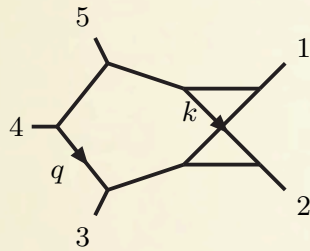
# PENTABOX INTEGRAND DECOMPOSITION

$$\begin{aligned}
 N(q, k) &= \Delta_{12345678}(q, k) + \\
 &\quad + \Delta_{1234568}(q, k)D_7 + \Delta_{1234578}(q, k)D_6 + \\
 &\quad + \Delta_{1234678}(q, k)D_5 + \Delta_{1235678}(q, k)D_4 = \\
 &= c_{12345678,0} + c_{12345678,1} (q \cdot p_1) + \\
 &\quad + c_{1234568,0}D_7 + c_{1234578,0}D_6 + \\
 &\quad + c_{1234678,0}D_5 + c_{1235678,0}D_4 ,
 \end{aligned}$$



# PENTACROSS INTEGRAND DECOMPOSITION

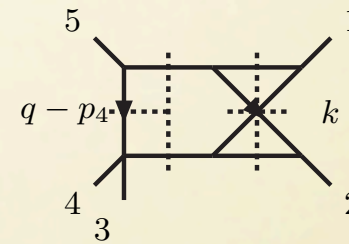
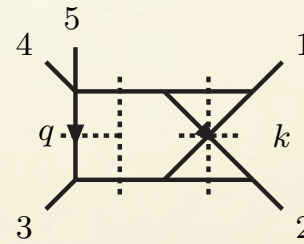
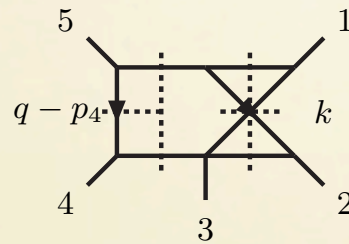
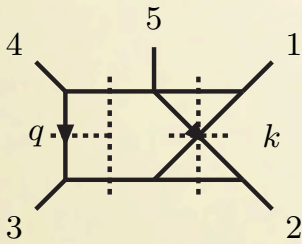
$$\begin{aligned}
 D_1 &= k^2 \\
 D_2 &= (k + p_2)^2 \\
 D_3 &= (k + q - p_4 - p_5)^2 \\
 D_4 &= q^2 \\
 D_5 &= (q + p_3)^2 \\
 D_6 &= (q - p_4)^2 \\
 D_7 &= (q - p_4 - p_5)^2 \\
 D_8 &= (q + k + p_2 + p_3)^2 .
 \end{aligned}$$



$$N(q, k) = 2q \cdot v + \alpha \quad \text{Carrasco \& Johansson (2011)}$$

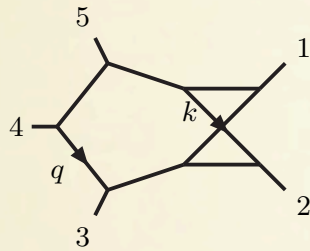
$$v^\mu = \frac{1}{4} \left( \gamma_{12}(p_1^\mu - p_2^\mu) + \gamma_{23}(p_2^\mu - p_3^\mu) + 2\gamma_{45}(p_4^\mu - p_5^\mu) + \gamma_{13}(p_1^\mu - p_3^\mu) \right)$$

$$\alpha = \frac{1}{4} \left( 2\gamma_{12}(s_{45} - s_{12}) + \gamma_{23}(s_{45} + 3s_{12} - s_{13}) + 2\gamma_{45}(s_{14} - s_{15}) + \gamma_{13}(s_{12} + s_{45} - s_{13}) \right)$$

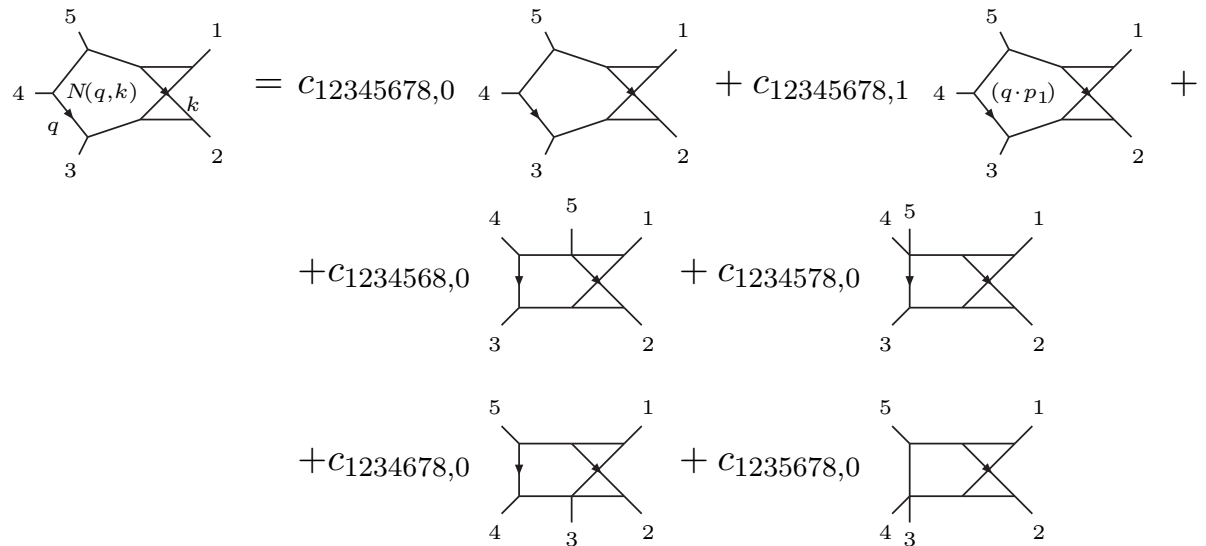


# PENTACROSS INTEGRAND DECOMPOSITION

$$\begin{aligned}
 D_1 &= k^2 \\
 D_2 &= (k + p_2)^2 \\
 D_3 &= (k + q - p_4 - p_5)^2 \\
 D_4 &= q^2 \\
 D_5 &= (q + p_3)^2 \\
 D_6 &= (q - p_4)^2 \\
 D_7 &= (q - p_4 - p_5)^2 \\
 D_8 &= (q + k + p_2 + p_3)^2 .
 \end{aligned}$$



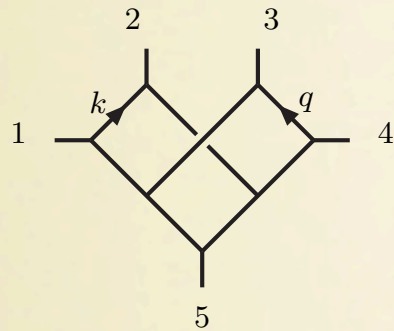
$$\begin{aligned}
 N(q, k) &= \Delta_{12345678}(q, k) + \\
 &+ \Delta_{1234568}(q, k)D_7 + \Delta_{1234578}(q, k)D_6 + \\
 &+ \Delta_{1234678}(q, k)D_5 + \Delta_{1235678}(q, k)D_4 = \\
 &= c_{12345678,0} + c_{12345678,1}(q \cdot p_1) + \\
 &+ c_{1234568,0}D_7 + c_{1234578,0}D_6 + \\
 &+ c_{1234678,0}D_5 + c_{1235678,0}D_4 ,
 \end{aligned}$$



The coefficients are the same of the planar case.



# THE LAST CONTRIBUTION TO THE 5-POINT N=4 SYM



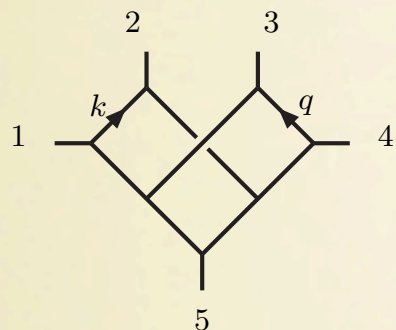
$$\begin{aligned}D_1 &= k^2 \\D_2 &= (k - p_1)^2 \\D_3 &= (k + p_2)^2 \\D_4 &= q^2 \\D_5 &= (q + p_3)^2 \\D_6 &= (q - p_4)^2 \\D_7 &= (q - k + p_1 + p_3)^2 \\D_8 &= (q - k - p_2 - p_4)^2\end{aligned}$$

$N(q,k)$  is *linear* in the loop momenta

Carrasco & Johansson (2011)

# 5-POINT 8FOLD-CUT

$D_1 = \dots = D_8 = 0$       8 solutions



$$\Delta_{12345678}(q, k) = \text{Res}_{12345678} \left\{ \mathcal{N}_{1\dots 8}(q, k) \right\} .$$

$$q^\mu = \sum_{i=1}^4 y_i \tau_i^\mu, \quad k^\mu = \sum_{i=1}^4 x_i e_i^\mu .$$

$$e_1 = p_1, \quad e_2 = p_2, \quad \tau_1 = p_3, \quad \tau_2 = p_4$$

The residue contains 8 monomials

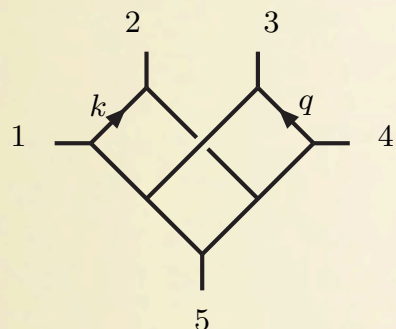
$$\{1, x_4, y_3, y_3^2, y_4, x_4 y_4, y_4^2, y_4^3\}$$

[Maximum Cut Thm]

generic residue

# 5-POINT 8FOLD-CUT

$$D_1 = \dots = D_8 = 0 \quad 8 \text{ solutions}$$



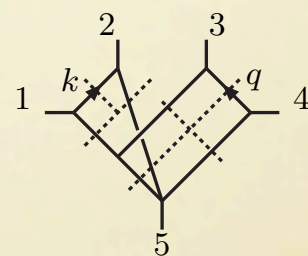
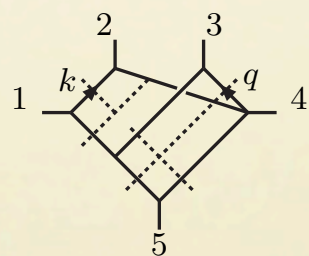
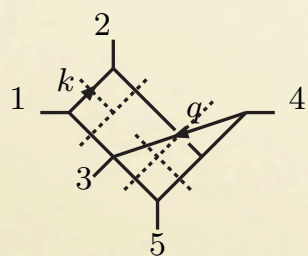
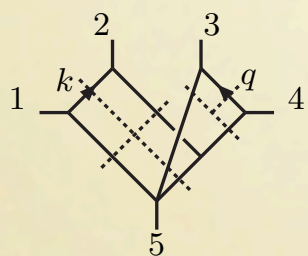
$$\Delta_{12345678}(q, k) = \text{Res}_{12345678} \left\{ \mathcal{N}_{1\dots 8}(q, k) \right\} .$$

$$q^\mu = \sum_{i=1}^4 y_i \tau_i^\mu, \quad k^\mu = \sum_{i=1}^4 x_i e_i^\mu .$$

$$e_1 = p_1, \quad e_2 = p_2, \quad \tau_1 = p_3, \quad \tau_2 = p_4$$

The residue contains 8 monomials  $\{1, x_4, y_3, y_3^2, y_4, x_4 y_4, y_4^2, y_4^3\}$

... FURTHER REDUCTION ...



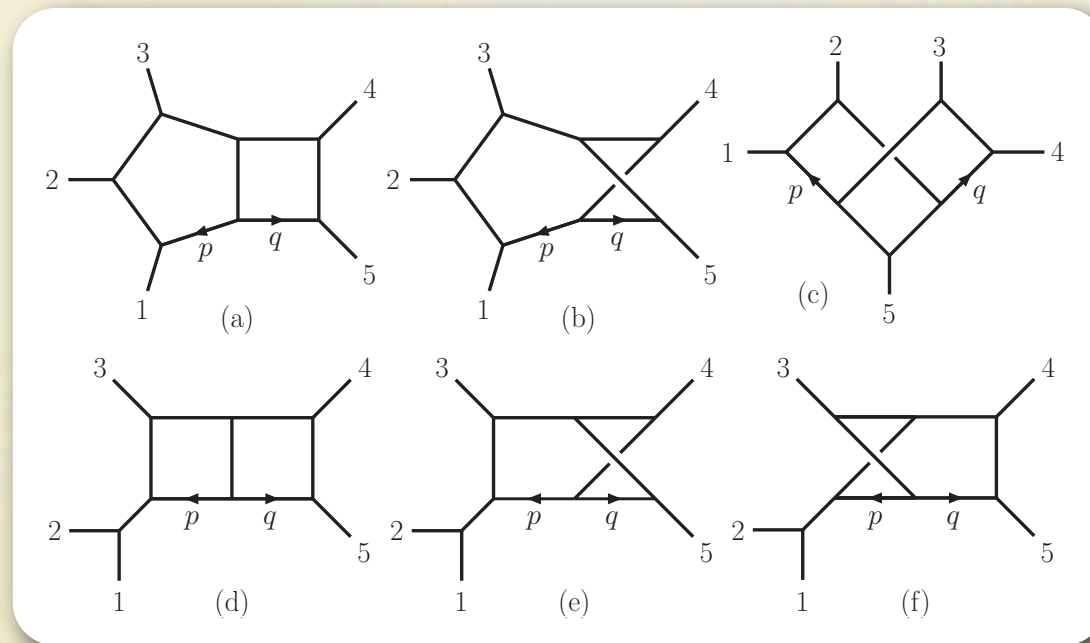
COMPLETE DECOMPOSITION

Global (N=N)-test fulfilled!

# 2-LOOP 5-POINT AMPLITUDES IN N=8 SUGRA

Same topologies as in the N=4 SYM, but  $N(q,k)$  is *quadratic* in the loop momenta

Carrasco & Johansson (2011)



The integrand reduction is analogous to the N=4 SYM case, involving the same cuts and residues.

Due to one extra power of loop momenta, the reduction involves also **6-denominator diagrams**: in the corresponding residues, the constant term is the only non-vanishing coefficient.

# CONCLUSIONS

- A unique mathematical framework for Amplitudes at any order in Perturbation Theory
  - one ingredient: Feynman denominator
  - one operation: *partial fractioning*
- Multivariate Polynomial Division/Groebner-basis generates the **residue** at an arbitrary cut
  - the general expression for the factorized amplitude
- Residues' **classification** complementary to Landau's singularity classification
- byproduct: the Maximum-cut Theorem
- *Recursive* generation of the *Integrand-decomposition Formula* @ any loop
- Amplitude decomposition from the shape of **residues**
  - ISP's determine a (non-minimal) MI-set
  - application: planar and non-planar 2-loop 5-point N=4 SYM and N=8 SuGra (low-rank numerators w.r.t. QCD)

to appear soon

Let me thank my collaborators who, directly or indirectly, have provided the material for this lectures:


- Uli Schubert
- Hans van Deurzen
- Tiziano Peraro
- Gionata Luisoni
- Edoardo Mirabella
- Giovanni Ossola

**EXTRA SLIDES**

## weak Nullstellensatz Theorem

**Theorem 1.2.3** (Weak Hilbert Nullstellensatz). *If  $k$  is algebraically closed, then  $V(S) = \emptyset$  iff there exists  $f_1 \dots f_N \in S$  and  $g_1 \dots g_N \in k[x_1, \dots, x_n]$  such that  $\sum f_i g_i = 1$*

The German word nullstellensatz could be translated as “zero set theorem”. The Weak Nullstellensatz can be rephrased as  $V(S) = \emptyset$  iff  $\langle S \rangle = (1)$ . Since this result is central to much of what follows, **we will assume that  $k$  is alge-**

 **Radical Ideal** Given an ideal  $\mathcal{J}$ , the *radical* of  $\mathcal{J}$  is  $\sqrt{\mathcal{J}} \equiv \{f \in P[\mathbf{z}] : \exists s \in \mathbb{N}, f^s \in \mathcal{J}\}$ .  
 $\mathcal{J}$  is radical iff  $\mathcal{J} = \sqrt{\mathcal{J}}$ .

## Finiteness Theorem

The following theorem bounds the number of points in  $\mathbf{V}(I)$  whenever  $I$  is zero-dimensional.

**Theorem 3-4.** *Let  $I$  be a zero-dimensional ideal in  $\mathbb{C}[x_1, \dots, x_n]$ . Then the number of points in  $\mathbf{V}(I)$  is at most  $\dim_{\mathbb{C}}(A)$ . Equality occurs if and only if  $I$  is a radical ideal.*

## Shape Lemma

since  $p_{red}(x)/x$  is a cubic polynomial in  $x^2$ .

If  $I$  is a zero-dimensional radical ideal in  $S = \mathbb{Q}[x_1, \dots, x_n]$  then, possibly after a linear change of variables, the ring  $S/I$  is always isomorphic to the univariate quotient ring  $\mathbb{Q}[x_i]/(I \cap \mathbb{Q}[x_i])$ . This is the content of the following result.

**PROPOSITION 2.3. (Shape Lemma)** *Let  $I$  be a zero-dimensional radical ideal in  $\mathbb{Q}[x_1, \dots, x_n]$  such that all  $d$  complex roots of  $I$  have distinct  $x_n$ -coordinates. Then the reduced Gröbner basis of  $I$  in the lexicographic term order has the shape*

$$\mathcal{G} = \{x_1 - q_1(x_n), x_2 - q_2(x_n), \dots, x_{n-1} - q_{n-1}(x_n), r(x_n)\}$$

where  $r$  is a polynomial of degree  $d$  and the  $q_i$  are polynomials of degree  $\leq d - 1$ .

For polynomial systems of moderate size, Singular is really doing the computing.



# MCT: PROOF (PART 1)

*Proof.* Let us parametrize the propagators using  $4\ell$  variables  $\mathbf{z} = (z_1, \dots, z_{4\ell})$ . In this parametrization, the solutions of the maximum-cut read,

$$\mathbf{z}^{(i)} = \left( z_1^{(i)}, \dots, z_{4\ell}^{(i)} \right), \text{ with } i = 1, \dots, n_s .$$

Let  $\mathcal{J}_{i_1 \dots i_{4\ell}}$  be the ideal generated by the on-shell denominators,  $\mathcal{J}_{i_1 \dots i_{4\ell}} = \langle D_{i_1}, \dots, D_{i_{4\ell}} \rangle$ . According to the assumptions, the number  $n_s$  of the solutions is finite, and each of them has multiplicity one, therefore  $\mathcal{J}_{i_1 \dots i_{4\ell}}$  is zero-dimensional and radical <sup>1</sup>, In this case, the *Finiteness Theorem* ensures that the remainder of the division of any polynomial modulo  $\mathcal{J}_{i_1 \dots i_{4\ell}}$  can be parametrised exactly by  $n_s$  coefficients.

## MCT: PROOF (PART 2)

Moreover, up to a linear coordinate change, we can assume that all the solutions of the system have distinct first coordinate  $z_1$ , i.e.  $z_1^{(i)} \neq z_1^{(j)} \forall i \neq j$ . We observe that  $\mathcal{J}_{i_1 \dots i_{4\ell}}$  and  $z_1$  are in the Shape Lemma position therefore a Gröbner basis for the lexicographic order  $z_1 < z_2 < \dots < z_n$  is  $\mathcal{G}_{i_1 \dots i_{4\ell}} = \{g_1, \dots, g_{4\ell}\}$ , in the form

$$\begin{cases} g_1(\mathbf{z}) = f_1(z_1) \\ g_2(\mathbf{z}) = z_2 - f_2(z_1) \\ \vdots \\ g_{4\ell}(\mathbf{z}) = z_{4\ell} - f_{4\ell}(z_1) . \end{cases}$$

The functions  $f_i$  are univariate polynomials in  $z_1$ . In particular  $f_1$  is a rank- $n_s$  square-free polynomial

$$f_1(z_1) = \prod_{i=1}^{n_s} (z_1 - z_1^{(i)}) ,$$

i.e. it does not exhibit repeated roots. The multivariate division of  $\mathcal{N}_{i_1 \dots i_{4\ell}}$  modulo  $\mathcal{G}_{i_1 \dots i_{4\ell}}$  leaves a remainder  $\Delta_{i_1 \dots i_{4\ell}}$  which is a univariate polynomial in  $z_1$  of degree  $(n_s - 1)$  in accordance with the *Finiteness Theorem*.  $\square$

# QCD recursion relations from the largest time equation Vaman, Yao (2005)

The factorization procedure is to cut these  $q_i$  successively by shifting them by  $z\eta$ . The on-shell conditions will give us a set of solutions, points in the complex plane, namely  $z_i = \frac{q_i^2 + m_i^2}{2\eta \cdot q_i}$ ,

The identity which we want to establish is

$$\begin{aligned} \frac{1}{q_1^2 + m_1^2} \frac{1}{q_2^2 + m_2^2} \cdots \frac{1}{q_{n-1}^2 + m_{n-1}^2} &= \frac{1}{q_1^2 + m_1^2} \frac{1}{(q_2 - z_1\eta)^2 + m_2^2} \cdots \frac{1}{(q_{n-1} - z_1\eta)^2 + m_{n-1}^2} \\ &+ \frac{1}{(q_1 - z_2\eta)^2 + m_1^2} \frac{1}{q_2^2 + m_2^2} \cdots \frac{1}{(q_{n-1} - z_2\eta)^2 + m_{n-1}^2} \\ &+ \cdots \cdots \cdots \\ &+ \frac{1}{(q_1 - z_{n-1}\eta)^2 + m_1^2} \cdots \frac{1}{(q_{n-2} - z_{n-1}\eta)^2 + m_{n-2}^2} \frac{1}{q_{n-1}^2 + m_{n-1}^2} \end{aligned}$$

making cuts, we have

$$\begin{aligned} \bar{q}_i^2 + m_i^2 = 0 &\rightarrow q_i^2 + m_i^2 = 2z_i\eta \cdot q_i & (q_i - z_j\eta)^2 + m_i^2 &= q_i^2 + m_i^2 - 2z_j\eta \cdot q_i \\ & & (q_i - z_j\eta)^2 + m_i^2 &= 2\eta \cdot q_i(z_i - z_j). \end{aligned}$$

Putting these together, we see the identity holds if one can show

$$\begin{aligned} \frac{(-1)^n}{z_1 z_2 \cdots z_{n-1}} &= \frac{1}{z_1(z_1 - z_2)(z_1 - z_3) \cdots (z_1 - z_{n-1})} \\ &+ \frac{1}{(z_2 - z_1)z_2(z_2 - z_3) \cdots (z_2 - z_{n-1})} \\ &\cdots \cdots \cdots \\ &+ \frac{1}{(z_{n-1} - z_1)(z_{n-1} - z_2) \cdots (z_{n-1} - z_{n-2})z_{n-1}}. \end{aligned} \tag{7.9}$$

This is so, because (7.9) is just a formula of partial fractioning, or it is just a statement that the integral

$$\int \frac{dz}{z(z - z_1)(z - z_2) \cdots (z - z_{n-1})} = 0$$

for a complex variable  $z$  over a contour which encloses all the poles.

### 2.2.2 Quintuple cut

The residue of the quintuple-cut,  $\bar{D}_i = \dots = \bar{D}_m = 0$ , defined as,

$$\Delta_{ijklm}(\bar{q}) = \text{Res}_{ijklm} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} \right\} = c_{5,0}^{(ijklm)} \mu^2 .$$

### 2.2.2 Quintuple cut

The residue of the quintuple-cut,  $\bar{D}_i = \dots = \bar{D}_m = 0$ , defined as,

$$\Delta_{ijklm}(\bar{q}) = \text{Res}_{ijklm} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} \right\} = c_{5,0}^{(ijklm)} \mu^2 .$$

### 2.2.3 Quadruple cut

The residue of the quadruple-cut,  $\bar{D}_i = \dots = \bar{D}_\ell = 0$ , defined as,

$$\begin{aligned} \Delta_{ijkl}(\bar{q}) &= \text{Res}_{ijkl} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i << m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} \right\} \\ &= c_{4,0}^{(ijkl)} + c_{4,2}^{(ijkl)} \mu^2 + c_{4,4}^{(ijkl)} \mu^4 - \left( c_{4,1}^{(ijkl)} + c_{4,3}^{(ijkl)} \mu^2 \right) \left[ (K_3 \cdot e_4) x_4 - (K_3 \cdot e_3) x_3 \right] (e_1 \cdot e_2) , \end{aligned}$$

### 2.2.2 Quintuple cut

The residue of the quintuple-cut,  $\bar{D}_i = \dots = \bar{D}_m = 0$ , defined as,

$$\Delta_{ijklm}(\bar{q}) = \text{Res}_{ijklm} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} \right\} = c_{5,0}^{(ijklm)} \mu^2 .$$

### 2.2.3 Quadruple cut

The residue of the quadruple-cut,  $\bar{D}_i = \dots = \bar{D}_\ell = 0$ , defined as,

$$\Delta_{ijkl}(\bar{q}) = \text{Res}_{ijkl} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i \ll m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} \right\} = c_{4,0}^{(ijkl)} + c_{4,2}^{(ijkl)} \mu^2 + c_{4,4}^{(ijkl)} \mu^4 - \left( c_{4,1}^{(ijkl)} + c_{4,3}^{(ijkl)} \mu^2 \right) \left[ (K_3 \cdot e_4) x_4 - (K_3 \cdot e_3) x_3 \right] (e_1 \cdot e_2),$$

### 2.2.4 Triple cut

The residue of the triple-cut,  $\bar{D}_i = \bar{D}_j = \bar{D}_k = 0$ , defined as,

$$\begin{aligned} \Delta_{ijk}(\bar{q}) &= \text{Res}_{ijk} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i \ll m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} - \sum_{i \ll \ell}^{n-1} \frac{\Delta_{ijkl}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell} \right\} \\ &= c_{3,0}^{(ijk)} + c_{3,7}^{(ijk)} \mu^2 - \left( (c_{3,1}^{(ijk)} + c_{3,8}^{(ijk)} \mu^2) x_4 + (c_{3,4}^{(ijk)} + c_{3,9}^{(ijk)} \mu^2) x_3 \right) (e_1 \cdot e_2) \\ &\quad + \left( c_{3,2}^{(ijk)} x_4^2 + c_{3,5}^{(ijk)} x_3^2 \right) (e_1 \cdot e_2)^2 - \left( c_{3,3}^{(ijk)} x_4^3 + c_{3,6}^{(ijk)} x_3^3 \right) (e_1 \cdot e_2)^3 . \end{aligned}$$

### 2.2.2 Quintuple cut

The residue of the quintuple-cut,  $\bar{D}_i = \dots = \bar{D}_m = 0$ , defined as,

$$\Delta_{ijklm}(\bar{q}) = \text{Res}_{ijklm} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} \right\} = c_{5,0}^{(ijklm)} \mu^2 .$$

### 2.2.3 Quadruple cut

The residue of the quadruple-cut,  $\bar{D}_i = \dots = \bar{D}_\ell = 0$ , defined as,

$$\Delta_{ijk\ell}(\bar{q}) = \text{Res}_{ijk\ell} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i << m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} \right\} = c_{4,0}^{(ijkl)} + c_{4,2}^{(ijkl)} \mu^2 + c_{4,4}^{(ijkl)} \mu^4 - \left( c_{4,1}^{(ijkl)} + c_{4,3}^{(ijkl)} \mu^2 \right) \left[ (K_3 \cdot e_4) x_4 - (K_3 \cdot e_3) x_3 \right] (e_1 \cdot e_2) ,$$

### 2.2.4 Triple cut

The residue of the triple-cut,  $\bar{D}_i = \bar{D}_j = \bar{D}_k = 0$ , defined as,

$$\begin{aligned} \Delta_{ijk}(\bar{q}) &= \text{Res}_{ijk} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i << m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} - \sum_{i << \ell}^{n-1} \frac{\Delta_{ijk\ell}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell} \right\} \\ &= c_{3,0}^{(ijk)} + c_{3,7}^{(ijk)} \mu^2 - \left( (c_{3,1}^{(ijk)} + c_{3,8}^{(ijk)} \mu^2) x_4 + (c_{3,4}^{(ijk)} + c_{3,9}^{(ijk)} \mu^2) x_3 \right) (e_1 \cdot e_2) + \left( c_{3,2}^{(ijk)} x_4^2 + c_{3,5}^{(ijk)} x_3^2 \right) (e_1 \cdot e_2)^2 - \left( c_{3,3}^{(ijk)} x_4^3 + c_{3,6}^{(ijk)} x_3^3 \right) (e_1 \cdot e_2)^3 . \end{aligned}$$

### 2.2.5 Double cut

The residue of the double-cut,  $\bar{D}_i = \bar{D}_j = 0$ , defined as,

$$\begin{aligned} \Delta_{ij}(\bar{q}) &= \text{Res}_{ij} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i << m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} - \sum_{i << \ell}^{n-1} \frac{\Delta_{ijk\ell}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell} - \sum_{i << k}^{n-1} \frac{\Delta_{ijk}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k} \right\} \\ &= c_{2,0}^{(ij)} + c_{2,9}^{(ij)} \mu^2 + \left( c_{2,1}^{(ij)} x_1 - c_{2,3}^{(ij)} x_4 - c_{2,5}^{(ij)} x_3 \right) (e_1 \cdot e_2) + \\ &\quad \left( c_{2,2}^{(ij)} x_1^2 + c_{2,4}^{(ij)} x_4^2 + c_{2,6}^{(ij)} x_3^2 - c_{2,7}^{(ij)} x_1 x_4 - c_{2,8}^{(ij)} x_1 x_3 \right) (e_1 \cdot e_2)^2 . \end{aligned}$$

### 2.2.2 Quintuple cut

The residue of the quintuple-cut,  $\bar{D}_i = \dots = \bar{D}_m = 0$ , defined as,

$$\Delta_{ijklm}(\bar{q}) = \text{Res}_{ijklm} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} \right\} = c_{5,0}^{(ijklm)} \mu^2 .$$

### 2.2.3 Quadruple cut

The residue of the quadruple-cut,  $\bar{D}_i = \dots = \bar{D}_\ell = 0$ , defined as,

$$\Delta_{ijkl\ell}(\bar{q}) = \text{Res}_{ijkl\ell} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i \ll m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} \right\} = c_{4,0}^{(ijkl)} + c_{4,2}^{(ijkl)} \mu^2 + c_{4,4}^{(ijkl)} \mu^4 - \left( c_{4,1}^{(ijkl)} + c_{4,3}^{(ijkl)} \mu^2 \right) \left[ (K_3 \cdot e_4) x_4 - (K_3 \cdot e_3) x_3 \right] (e_1 \cdot e_2) ,$$

### 2.2.4 Triple cut

The residue of the triple-cut,  $\bar{D}_i = \bar{D}_j = \bar{D}_k = 0$ , defined as,

$$\begin{aligned} \Delta_{ijk}(\bar{q}) &= \text{Res}_{ijk} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i \ll m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} - \sum_{i \ll \ell}^{n-1} \frac{\Delta_{ijk\ell}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell} \right\} \\ &= c_{3,0}^{(ijk)} + c_{3,7}^{(ijk)} \mu^2 - \left( (c_{3,1}^{(ijk)} + c_{3,8}^{(ijk)} \mu^2) x_4 + (c_{3,4}^{(ijk)} + c_{3,9}^{(ijk)} \mu^2) x_3 \right) (e_1 \cdot e_2) + \left( c_{3,2}^{(ijk)} x_4^2 + c_{3,5}^{(ijk)} x_3^2 \right) (e_1 \cdot e_2)^2 - \left( c_{3,3}^{(ijk)} x_4^3 + c_{3,6}^{(ijk)} x_3^3 \right) (e_1 \cdot e_2)^3 . \end{aligned}$$

### 2.2.5 Double cut

#### 2.2.6 Single cut

The residue of the single-cut,  $\bar{D}_i = 0$ , defined as,

$$\begin{aligned} \Delta_i(\bar{q}) &= \text{Res}_i \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i \ll m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} - \sum_{i \ll \ell}^{n-1} \frac{\Delta_{ijk\ell}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell} - \sum_{i \ll k}^{n-1} \frac{\Delta_{ijk}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k} - \sum_{i < j}^{n-1} \frac{\Delta_{ij}(\bar{q})}{\bar{D}_i \bar{D}_j} \right\} \\ &= c_{1,0}^{(i)} + \left( c_{1,1}^{(i)} x_2 + c_{1,2}^{(i)} x_1 - c_{1,3}^{(i)} x_4 - c_{1,4}^{(i)} x_3 \right) (e_1 \cdot e_2) . \end{aligned}$$



### 2.2.2 Quintuple cut

The residue of the quintuple-cut,  $\bar{D}_i = \dots = \bar{D}_m = 0$ , defined as,

$$\Delta_{ijklm}(\bar{q}) = \text{Res}_{ijklm} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} \right\} = c_{5,0}^{(ijklm)} \mu^2 .$$

### 2.2.3 Quadruple cut

The residue of the quadruple-cut,  $\bar{D}_i = \dots = \bar{D}_\ell = 0$ , defined as,

$$\Delta_{ijkl}(\bar{q}) = \text{Res}_{ijkl} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i < \ell < m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} \right\} = c_{4,0}^{(ijkl)} + c_{4,2}^{(ijkl)} \mu^2 + c_{4,4}^{(ijkl)} \mu^4 - \left( c_{4,1}^{(ijkl)} + c_{4,3}^{(ijkl)} \mu^2 \right) \left[ (K_3 \cdot e_4) x_4 - (K_3 \cdot e_3) x_3 \right] (e_1 \cdot e_2) ,$$

### 2.2.4 Triple cut

The residue of the triple-cut,  $\bar{D}_i = \bar{D}_j = \bar{D}_k = 0$ , defined as,

$$\Delta_{ijk}(\bar{q}) = \text{Res}_{ijk} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i < \ell < m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} - \sum_{i < \ell}^{n-1} \frac{\Delta_{ijkl}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell} \right\}$$

**SAMURAI** Ossola, Reiter, Tramontano, & P.M. (2010)

**SCATTERING AMPLITUDES FROM UNITARITY-BASED REDUCTION ALGORITHM AT THE INTEGRAND-LEVEL**

$$= c_{2,0}^{(ij)} + c_{2,9}^{(ij)} \mu^2 + \left( c_{2,1}^{(ij)} x_1 - c_{2,3}^{(ij)} x_4 - c_{2,5}^{(ij)} x_3 \right) (e_1 \cdot e_2) + \left( c_{2,2}^{(ij)} x_1^2 + c_{2,4}^{(ij)} x_4^2 + c_{2,6}^{(ij)} x_3^2 - c_{2,7}^{(ij)} x_1 x_4 - c_{2,8}^{(ij)} x_1 x_3 \right) (e_1 \cdot e_2)^2 .$$

### 2.2.6 Single cut

The residue of the single-cut,  $\bar{D}_i = 0$ , defined as,

$$\Delta_i(\bar{q}) = \text{Res}_i \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i < \ell < m}^{n-1} \frac{\Delta_{ijklm}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} - \sum_{i < \ell}^{n-1} \frac{\Delta_{ijkl}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell} - \sum_{i < k}^{n-1} \frac{\Delta_{ijk}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k} - \sum_{i < j}^{n-1} \frac{\Delta_{ij}(\bar{q})}{\bar{D}_i \bar{D}_j} \right\}$$

$$= c_{1,0}^{(i)} + \left( c_{1,1}^{(i)} x_2 + c_{1,2}^{(i)} x_1 - c_{1,3}^{(i)} x_4 - c_{1,4}^{(i)} x_3 \right) (e_1 \cdot e_2) .$$

## 2.3 Amplitude and master integrals

$$\begin{aligned}
\mathcal{A}_n = & \sum_{i < j < k < \ell}^{n-1} \left\{ c_{4,0}^{(ijkl)} I_{ijkl}^{(d)} + \frac{(d-2)(d-4)}{4} c_{4,4}^{(ijkl)} I_{ijkl}^{(d+4)} \right\} \\
& + \sum_{i < j < k}^{n-1} \left\{ c_{3,0}^{(ijk)} I_{ijk}^{(d)} - \frac{(d-4)}{2} c_{3,7}^{(ijk)} I_{ijk}^{(d+2)} \right\} \\
& + \sum_{i < j}^{n-1} \left\{ c_{2,0}^{(ij)} I_{ij}^{(d)} + c_{2,1}^{(ij)} J_{ij}^{(d)} + c_{2,2}^{(ij)} K_{ij}^{(d)} - \frac{(d-4)}{2} c_{2,9}^{(ij)} I_{ij}^{(d+2)} \right\} \\
& + \sum_i^{n-1} c_{1,0}^{(i)} I_i^{(d)},
\end{aligned}$$

$$\begin{aligned}
\int d^d \bar{q} \frac{\mu^4}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell} &= \frac{(d-2)(d-4)}{4} I_{ijkl}^{(d+4)}, \\
\int d^d \bar{q} \frac{\mu^2}{\bar{D}_i \bar{D}_j \bar{D}_k} &= -\frac{(d-4)}{2} I_{ijk}^{(d+2)}, \\
\int d^d \bar{q} \frac{\mu^2}{\bar{D}_i \bar{D}_j} &= -\frac{(d-4)}{2} I_{ij}^{(d+2)}, \\
\int d^d \bar{q} \frac{\bar{q} \cdot e_2}{\bar{D}_i \bar{D}_j} &= J_{ij}^{(d)}, \\
\int d^d \bar{q} \frac{(\bar{q} \cdot e_2)^2}{\bar{D}_i \bar{D}_j} &= K_{ij}^{(d)}.
\end{aligned}$$

$$\begin{aligned}
\int \frac{d^4 p}{(2\pi)^4} \int \frac{d^{-2\epsilon} \mu}{(2\pi)^{-2\epsilon}} (\mu^2)^r f(p^\alpha, \mu^2) &= \int \frac{d^4 p}{(2\pi)^4} \int d\Omega_{-1-2\epsilon} \int_0^\infty \frac{d\mu^2}{2(2\pi)^{-2\epsilon}} (\mu^2)^{-1-\epsilon+r} f(p^\alpha, \mu^2) \\
&= \frac{(2\pi)^{2r} \int d\Omega_{-1-2\epsilon}}{\int d\Omega_{2r-1-2\epsilon}} \int \frac{d^4 p}{(2\pi)^4} \int \frac{d^{2r-2\epsilon} \mu}{(2\pi)^{2r-2\epsilon}} f(p^\alpha, \mu^2) \\
&= -\epsilon(1-\epsilon)(2-\epsilon) \cdots (r-1-\epsilon) (4\pi)^r \int \frac{d^{4+2r-2\epsilon} P}{(2\pi)^{4+2r-2\epsilon}} f(p^\alpha, \mu^2),
\end{aligned}$$

## ■ 2-D Minkowski product

```
MP[a_, b_] := a[[1]] * b[[1]] - a[[2]] * b[[2]] ;
```

## ■ Kinematics

```
N[P11] = 0.3 * Sqrt[11] // N;  
N[P12] = 0.5 * Sqrt[13] // N;  
N[P21] = -0.1 * Sqrt[17] // N;  
N[P22] = 0.7 * Sqrt[19] // N;
```

```
P1 = {P11, P12};  
P2 = {P21, P22};  
P3 = -P1 - P2;
```

```
MP[P1, P1] // N  
MP[P2, P2] // N  
MP[P3, P3] // N
```

-2.26

-9.14

-23.2219

## ■ orthogonal momentum: PT.P1=0

```
PT = {P12, P11}
```

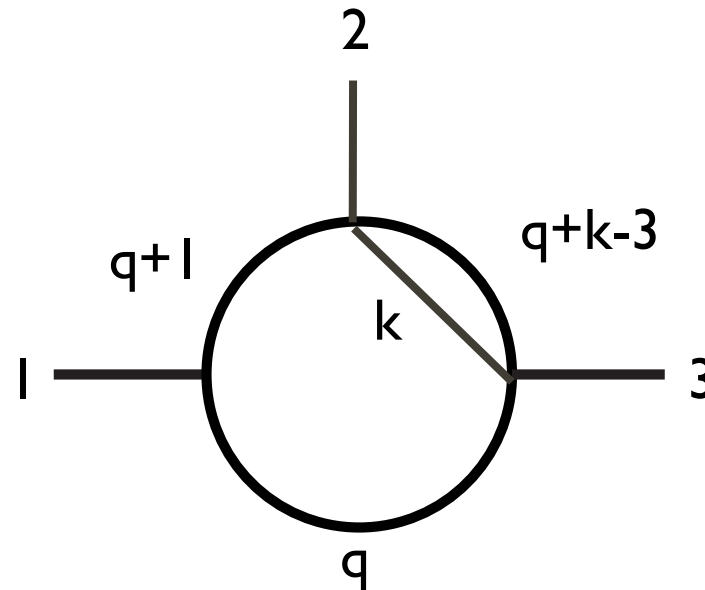
```
{P12, P11}
```

```
MP[PT, PT] // N
```

2.26

```
MP[P1, PT] // N
```

0.



2D Maximum Cut

## ■ Loop momenta decomposition

$$\mathbf{q} = x_1 * \mathbf{P}_1 + x_2 * \mathbf{P}_T;$$

$$\mathbf{k} = y_1 * \mathbf{P}_1 + y_2 * \mathbf{P}_T;$$

## ■ scalar products vs components

**MP[q, P1] // N // Expand // Chop**

$$-2.26 x_1$$

**MP[q, PT] // N // Expand // Chop**

$$2.26 x_2$$

**MP[k, P1] // N // Expand // Chop**

$$-2.26 y_1$$

**MP[k, PT] // N // Expand // Chop**

$$2.26 y_2$$

**Vars = {x1, x2, y1, y2};**

## ■ Denominators

$$D_1 = \text{MP}[\mathbf{q}, \mathbf{q}];$$

$$D_2 = \text{MP}[\mathbf{q} + \mathbf{P}_1, \mathbf{q} + \mathbf{P}_1];$$

$$D_3 = \text{MP}[\mathbf{q} + \mathbf{k} - \mathbf{P}_3, \mathbf{q} + \mathbf{k} - \mathbf{P}_3];$$

$$D_4 = \text{MP}[\mathbf{k}, \mathbf{k}];$$

**Dens = {D1, D2, D3, D4} // N // Expand // Chop;**

**For[i = 1, i ≤ 4, i++,**

**Print[Dens[[i]]];**

**]**

$$-2.26 x_1^2 + 2.26 x_2^2$$

$$-2.26 - 4.52 x_1 - 2.26 x_1^2 + 2.26 x_2^2$$

$$-23.2219 - 16.3419 x_1 - 2.26 x_1^2 - 7.55848 x_2 + 2.26 x_2^2 -$$

$$16.3419 y_1 - 4.52 x_1 y_1 - 2.26 y_1^2 - 7.55848 y_2 + 4.52 x_2 y_2 + 2.26 y_2^2$$

$$-2.26 y_1^2 + 2.26 y_2^2$$

## ■ On-shell Cut-conditions

```
Cut = Table[Dens[[i]] == 0, {i, 1, 4}]
```

$$\{-2.26 x_1^2 + 2.26 x_2^2 = 0, -2.26 - 4.52 x_1 - 2.26 x_1^2 + 2.26 x_2^2 = 0, \\ -23.2219 - 16.3419 x_1 - 2.26 x_1^2 - 7.55848 x_2 + 2.26 x_2^2 - 16.3419 y_1 - 4.52 x_1 y_1 - \\ 2.26 y_1^2 - 7.55848 y_2 + 4.52 x_2 y_2 + 2.26 y_2^2 = 0, -2.26 y_1^2 + 2.26 y_2^2 = 0\}$$

## ■ Solutions

```
Sols = NSolve[Cut, Vars];
```

```
For[i = 1, i ≤ 4, i++,  
  Print[Sols[[i]]];  
]
```

```
{x1 → -0.5, x2 → -0.5, y1 → -2.64384, y2 → 2.64384}
```

```
{x1 → -0.5, x2 → 0.5, y1 → -2.14384, y2 → 2.14384}
```

```
{x1 → -0.5, x2 → 0.5, y1 → -0.971612, y2 → -0.971612}
```

```
{x1 → -0.5, x2 → -0.5, y1 → -0.471612, y2 → -0.471612}
```

```
Length[Sols]
```

```
4
```

## Rank-5 numerator

```

Num = (
  c[0] +
  Sum[c[i1] * Vars[[i1]], {i1, 1, Length[Vars]}] +
  Sum[c[i1, i2] * Vars[[i1]] * Vars[[i2]],
  {i1, 1, Length[Vars]}, {i2, i1, Length[Vars]}] +
  Sum[c[i1, i2, i3] * Vars[[i1]] * Vars[[i2]] * Vars[[i3]],
  {i1, 1, Length[Vars]}, {i2, i1, Length[Vars]}, {i3, i2, Length[Vars]}] +
  Sum[c[i1, i2, i3, i4] * Vars[[i1]] * Vars[[i2]] * Vars[[i3]] * Vars[[i4]],
  {i1, 1, Length[Vars]}, {i2, i1, Length[Vars]},
  {i3, i2, Length[Vars]}, {i4, i3, Length[Vars]}] +
  Sum[c[i1, i2, i3, i4, i5] * Vars[[i1]] * Vars[[i2]] * Vars[[i3]] *
  Vars[[i4]] * Vars[[i5]], {i1, 1, Length[Vars]}, {i2, i1, Length[Vars]},
  {i3, i2, Length[Vars]}, {i4, i3, Length[Vars]}, {i5, i4, Length[Vars]}]
)

```

```

c[0] + x1 c[1] + x2 c[2] + y1 c[3] + y2 c[4] + x12 c[1, 1] + x1 x2 c[1, 2] + x1 y1 c[1, 3] +
x1 y2 c[1, 4] + x22 c[2, 2] + x2 y1 c[2, 3] + x2 y2 c[2, 4] + y12 c[3, 3] + y1 y2 c[3, 4] +
y22 c[4, 4] + x13 c[1, 1, 1] + x12 x2 c[1, 1, 2] + x12 y1 c[1, 1, 3] + x12 y2 c[1, 1, 4] +
x1 x22 c[1, 2, 2] + x1 x2 y1 c[1, 2, 3] + x1 x2 y2 c[1, 2, 4] + x1 y12 c[1, 3, 3] +
x1 y1 y2 c[1, 3, 4] + x1 y22 c[1, 4, 4] + x23 c[2, 2, 2] + x22 y1 c[2, 2, 3] +
x22 y2 c[2, 2, 4] + x2 y12 c[2, 3, 3] + x2 y1 y2 c[2, 3, 4] + x2 y22 c[2, 4, 4] +
y13 c[3, 3, 3] + y12 y2 c[3, 3, 4] + y1 y22 c[3, 4, 4] + y23 c[4, 4, 4] + x14 c[1, 1, 1, 1] +
x13 x2 c[1, 1, 1, 2] + x13 y1 c[1, 1, 1, 3] + x13 y2 c[1, 1, 1, 4] + x12 x22 c[1, 1, 2, 2] +
x12 x2 y1 c[1, 1, 2, 3] + x12 x2 y2 c[1, 1, 2, 4] + x12 y12 c[1, 1, 3, 3] + x12 y1 y2 c[1, 1, 3, 4] +
x12 y22 c[1, 1, 4, 4] + x1 x23 c[1, 2, 2, 2] + x1 x22 y1 c[1, 2, 2, 3] + x1 x22 y2 c[1, 2, 2, 4] +
x1 x2 y12 c[1, 2, 3, 3] + x1 x2 y1 y2 c[1, 2, 3, 4] + x1 x2 y22 c[1, 2, 4, 4] +
x1 y13 c[1, 3, 3, 3] + x1 y12 y2 c[1, 3, 3, 4] + x1 y1 y22 c[1, 3, 4, 4] + x1 y23 c[1, 4, 4, 4] +
x24 c[2, 2, 2, 2] + x23 y1 c[2, 2, 2, 3] + x23 y2 c[2, 2, 2, 4] + x22 y12 c[2, 2, 3, 3] +
x22 y1 y2 c[2, 2, 3, 4] + x22 y22 c[2, 2, 4, 4] + x2 y13 c[2, 3, 3, 3] + x2 y12 y2 c[2, 3, 3, 4] +
x2 y1 y22 c[2, 3, 4, 4] + x2 y23 c[2, 4, 4, 4] + y14 c[3, 3, 3, 3] + y13 y2 c[3, 3, 3, 4] +
y12 y22 c[3, 3, 4, 4] + y1 y23 c[3, 4, 4, 4] + y24 c[4, 4, 4, 4] + x15 c[1, 1, 1, 1, 1] +
x14 x2 c[1, 1, 1, 1, 2] + x14 y1 c[1, 1, 1, 1, 3] + x14 y2 c[1, 1, 1, 1, 4] +
x13 x22 c[1, 1, 1, 2, 2] + x13 x2 y1 c[1, 1, 1, 2, 3] + x13 x2 y2 c[1, 1, 1, 2, 4] +
x13 y12 c[1, 1, 1, 3, 3] + x13 y1 y2 c[1, 1, 1, 3, 4] + x13 y22 c[1, 1, 1, 4, 4] +
x12 x23 c[1, 1, 2, 2, 2] + x12 x22 y1 c[1, 1, 2, 2, 3] + x12 x22 y2 c[1, 1, 2, 2, 4] +
x12 x2 y12 c[1, 1, 2, 3, 3] + x12 x2 y1 y2 c[1, 1, 2, 3, 4] + x12 x2 y22 c[1, 1, 2, 4, 4] +
x12 y13 c[1, 1, 3, 3, 3] + x12 y12 y2 c[1, 1, 3, 3, 4] + x12 y1 y22 c[1, 1, 3, 4, 4] +
x12 y23 c[1, 1, 4, 4, 4] + x1 x24 c[1, 2, 2, 2, 2] + x1 x23 y1 c[1, 2, 2, 2, 3] +
x1 x23 y2 c[1, 2, 2, 2, 4] + x1 x22 y12 c[1, 2, 2, 3, 3] + x1 x22 y1 y2 c[1, 2, 2, 3, 4] +
x1 x22 y22 c[1, 2, 2, 4, 4] + x1 x2 y13 c[1, 2, 3, 3, 3] + x1 x2 y12 y2 c[1, 2, 3, 3, 4] +
x1 x2 y1 y22 c[1, 2, 3, 4, 4] + x1 x2 y23 c[1, 2, 4, 4, 4] + x1 y14 c[1, 3, 3, 3, 3] +
x1 y13 y2 c[1, 3, 3, 3, 4] + x1 y12 y22 c[1, 3, 3, 4, 4] + x1 y1 y23 c[1, 3, 4, 4, 4] +
x1 y24 c[1, 4, 4, 4, 4] + x25 c[2, 2, 2, 2, 2] + x24 y1 c[2, 2, 2, 2, 3] + x24 y2 c[2, 2, 2, 2, 4] +
x23 y12 c[2, 2, 2, 3, 3] + x23 y1 y2 c[2, 2, 2, 3, 4] + x23 y22 c[2, 2, 2, 4, 4] +
x22 y13 c[2, 2, 3, 3, 3] + x22 y12 y2 c[2, 2, 3, 3, 4] + x22 y1 y22 c[2, 2, 3, 4, 4] +
x22 y23 c[2, 2, 4, 4, 4] + x2 y14 c[2, 3, 3, 3, 3] + x2 y13 y2 c[2, 3, 3, 3, 4] +
x2 y12 y22 c[2, 3, 3, 4, 4] + x2 y1 y23 c[2, 3, 4, 4, 4] + x2 y24 c[2, 4, 4, 4, 4] +
y15 c[3, 3, 3, 3, 3] + y14 y2 c[3, 3, 3, 3, 4] + y13 y22 c[3, 3, 3, 4, 4] +
y12 y23 c[3, 3, 4, 4, 4] + y1 y24 c[3, 4, 4, 4, 4] + y25 c[4, 4, 4, 4, 4]

```

Length[%]

## ■ Polynomial Division

```
GB = GroebnerBasis[Dens, Vars, MonomialOrder → DegreeLexicographic];
```

```
For[i = 1, i ≤ 4, i++,  
  Print[GB[[i]]];  
]
```

```
0.5 + 1. x1
```

```
1.66493 + 0.836115 x2 + 3.11545 y1 + 1. y22
```

```
-0.960748 - 0.983196 x2 - 1.67223 y1 + 1.14906 y2 + 1. y1 y2
```

```
1.66493 + 0.836115 x2 + 3.11545 y1 + 1. y12
```

```
remainder = PolynomialReduce[Num, GB, Vars, MonomialOrder → DegreeLexicographic][[2]];
```

```
Length[MonomialList[remainder, Vars]]
```

```
4
```

```
i = 1;
```

```
Collect[remainder, Vars, coeff[i++] &]
```

```
x2 coeff[1] + y1 coeff[2] + y2 coeff[3] + coeff[4]
```

## ■ Polynomial Division & Shape Lemma

```
GB = GroebnerBasis[Dens, Vars];
```

```
For[i = 1, i ≤ 4, i++,  
  Print[GB[[i]]];  
]
```

```
2.5972 + 5.98633 y2 - 0.7835 y22 - 3.34446 y23 + 1. y24
```

```
0.352643 + 1. y1 + 0.0586645 y2 + 0.605742 y22 - 0.113526 y23
```

```
0.677279 + 1. x2 - 0.21859 y2 - 1.06105 y22 + 0.423009 y23
```

```
0.5 + 1. x1
```

```
remainder = PolynomialReduce[Num, GB, Vars][[2]];
```

```
Length[MonomialList[remainder, Vars]]
```

```
4
```

```
i = 1;
```

```
Collect[remainder, Vars, coeff[i++] &]
```

```
y23 coeff[1] + y22 coeff[2] + y2 coeff[3] + coeff[4]
```